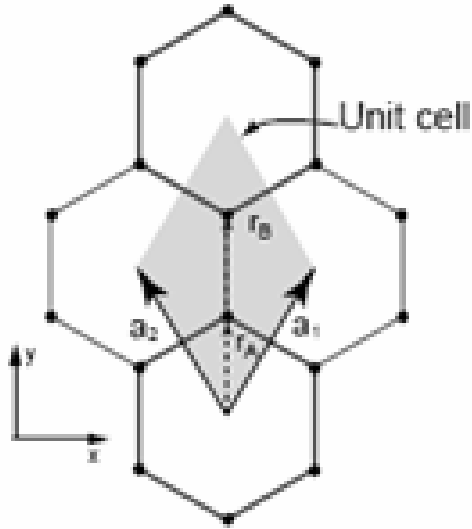


PRIBLIŽEK TESNE VEZI: GRAFEN

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Grafen ima 2D heksagonalno mrežo. Za vektorja Bravaisove mreže vzamemo



$$\vec{a}_1 = \left(\frac{a}{2}, \frac{\sqrt{3}}{2}a \right)$$

$$\vec{a}_2 = \left(-\frac{a}{2}, \frac{\sqrt{3}}{2}a \right)$$

$$\vec{R} = n_1\vec{a}_1 + n_2\vec{a}_2$$

Sedaj lahko izračunamo elektronsko strukturo s pomočjo približka tesne vezi. K strukturi prispeva le ena orbitala na posamezen atom, v posamezni osnovni celici pa imamo 2 atoma. Najprej zapišemo Blochovo funkcijo

$$\Phi_\alpha(\vec{k}, \vec{r}) = \sum_{\vec{R}} \exp(i\vec{k}\vec{R}) \phi_\alpha(\vec{r} - \vec{R})$$

kjer je ϕ_α valovna funkcija P_z orbitale.

Indeks α teče od 1 do n in predstavlja število orbital v osnovni celici. V našem primeru od 1 do 2, ker imamo 2 atoma.

Lastne funkcije elektronov sedaj lahko zapišemo kot linearno kombinacijo Blochovih funkcij za atoma \vec{r}_A in \vec{r}_B .

$$\Psi_j(\vec{k}, \vec{r}) = \sum_{\alpha=1}^2 c_{j,\alpha}(\vec{k}) \Phi_\alpha(\vec{k}, \vec{r})$$

Valovno funkcijo Ψ_j vstavimo v Schrödingerjevo enačbo. Koeficient $c_{j,\alpha}$ mora biti izbran tako, da ima Ψ_j najmanjšo energijo.

Tako pridemo do enačbe

$$\det[\kappa - \varepsilon_i \zeta] = 0$$

Z ε_i označimo lastne vrednosti energije v približku tesne vezi.
 κ, ζ sta 2×2 matriki

$$\begin{aligned}\kappa_{j,j}(\vec{k}) &= \langle \Phi_j | H | \Phi_{j'} \rangle \\ \zeta_{j,j}(\vec{k}) &= \langle \Phi_j | \Phi_{j'} \rangle\end{aligned}$$

Z indeksoma A in B smo označili atoma v osnovni celici

$$\begin{aligned}\kappa_{AA} &= \kappa_{BB} = \langle \Phi_A | H | \Phi_B \rangle \\ &= \sum_{\vec{R}, \vec{R}'} \exp(i\vec{k}(\vec{R} - \vec{R}')) \langle \phi_A(\vec{r} - \vec{R}') | H | \phi_A(\vec{r} - \vec{R}) \rangle \\ &= \sum_{\vec{R}=\vec{R}'} \langle \phi_A(\vec{r} - \vec{R}') | H | \phi_A(\vec{r} - \vec{R}') \rangle + \sum_{\vec{R} \neq \vec{R}'} \exp(i\vec{k}(\vec{R} - \vec{R}')) \langle \phi_A(\vec{r} - \vec{R}') | H | \phi_A(\vec{r} - \vec{R}) \rangle \\ &= \sum_{\vec{R}=\vec{R}'} \varepsilon_0 + \text{členi višjega reda} \approx \varepsilon_0 \\ \kappa_{AB} &= \kappa_{BA}^* = \sum_{\vec{R}, \vec{R}'} \exp(i\vec{k}(\vec{R} - \vec{R}')) \langle \phi_A(\vec{r} - \vec{R}') | H | \phi_B(\vec{r} - \vec{R}) \rangle \\ &= \sum_{\vec{R}=\vec{R}'} \langle \phi_A(\vec{r} - \vec{R}') | H | \phi_B(\vec{r} - \vec{R}') \rangle + \sum_{\vec{R}'=\vec{R}+\vec{a}_1} \exp(-i\vec{k}\vec{a}_1) \langle \phi_A(\vec{r} - \vec{R}') | H | \phi_B(\vec{r} - \vec{R}) \rangle \\ &\quad + \sum_{\vec{R}'=\vec{R}+\vec{a}_2} \exp(-i\vec{k}\vec{a}_2) \langle \phi_A(\vec{r} - \vec{R}') | H | \phi_B(\vec{r} - \vec{R}) \rangle + \text{členi višjega reda} \\ &= t(1 + \exp(-i\vec{k}\vec{a}_1) + \exp(-i\vec{k}\vec{a}_2)) = tg(\vec{k})\end{aligned}$$

Iz simetrije je vidno

$$\begin{aligned}&\langle \phi_A(\vec{r} - \vec{R}) | H | \phi_B(\vec{r} - \vec{R}) \rangle \\ &= \langle \phi_A(\vec{r} - (\vec{R} + \vec{a}_1)) | H | \phi_B(\vec{r} - \vec{R}) \rangle \\ &= \langle \phi_A(\vec{r} - (\vec{R} + \vec{a}_2)) | H | \phi_B(\vec{r} - \vec{R}) \rangle = t\end{aligned}$$

$$\begin{aligned}\zeta_{AA} &= \zeta_{BB} = 1 \\ \zeta_{AB} &= \zeta_{BA} = \sum_{\vec{R}, \vec{R}'} \exp(i\vec{k}(\vec{R} - \vec{R}')) \langle \Phi_A(\vec{r} - \vec{R}') | \Phi_B(\vec{r} - \vec{R}') \rangle\end{aligned}$$

$$\begin{aligned}
&= \sum_{\vec{R}=\vec{R}'} \langle \Phi_A(\vec{r}-\vec{R}') | \Phi_B(\vec{r}-\vec{R}') \rangle + \sum_{\vec{R}'-(\vec{R}+\vec{a}_1)} \exp(-i\vec{k}\vec{a}_1) \langle \phi_A | \phi_B \rangle \\
&+ \sum_{\vec{R}'-(\vec{R}+\vec{a}_2)} \exp(-i\vec{k}\vec{a}_2) \langle \phi_A | \phi_B \rangle + \text{členi višjega reda} = \tilde{t}g(\vec{k}) \\
&\tilde{t} = \langle \phi_A(\vec{r}-\vec{R}) | \phi_B(\vec{r}-\vec{R}) \rangle \\
&= \langle \phi_A(\vec{r}-(\vec{R}+\vec{a}_1)) | \phi_B(\vec{r}-\vec{R}) \rangle = \langle \phi_A(\vec{r}-(\vec{R}+\vec{a}_2)) | \phi_B(\vec{r}-\vec{R}) \rangle
\end{aligned}$$

$$\begin{aligned}
\kappa &= \begin{pmatrix} \varepsilon_0 & tg(\vec{k}) \\ tg(\vec{k}) & \varepsilon_0 \end{pmatrix} \quad \zeta = \begin{pmatrix} 1 & \tilde{t}g(\vec{k}) \\ \tilde{t}g(\vec{k}) & 1 \end{pmatrix} \\
\det[\kappa - \varepsilon_i \zeta] &= \begin{vmatrix} \varepsilon_0 - \varepsilon_i & (\tilde{t} - \varepsilon_i t)g(\vec{k}) \\ (\tilde{t} - \varepsilon_i t)g(\vec{k}) & \varepsilon_0 - \varepsilon_i \end{vmatrix} = 0 \\
\varepsilon_i(\vec{k}) &= \frac{\varepsilon_0 \pm |g(\vec{k})|t}{1 \pm |g(\vec{k})|\tilde{t}}
\end{aligned}$$

Člen $|g(\vec{k})|\tilde{t}$ zanemarim in ε_0 postavim na 0.

$$\varepsilon_i(\vec{k}) = \pm |g(\vec{k})|t$$

$$|g(\vec{k})| = \sqrt{1 + 4 \cos\left(\frac{\sqrt{3}k_y a}{2}\right) \cos\left(\frac{k_x a}{2}\right) + 4 \cos^2\left(\frac{k_x a}{2}\right)}$$

$$\varepsilon_i = 0 \implies |g(\vec{k})| = 0$$

$$\vec{k}_{1,2} = \left(\pm \frac{4\pi}{3a}, 0\right) \quad \vec{k}_{3,4,5,6} = \left(\pm \frac{2\pi}{3a}, \pm \frac{2\sqrt{3}}{3a}\pi\right)$$

Disperzija energije v bližini točke, kjer je $\varepsilon_i = 0$

$$\vec{k} = \vec{k}_0 + \delta\vec{k} \quad \delta\vec{k} = (\delta k_x, \delta k_y) \quad \delta\vec{k} \text{ majhen}$$

$$\begin{aligned}
\varepsilon(\vec{k}) &= \pm t |1 + \exp(-i(\vec{k}_0 + \delta\vec{k})\vec{a}_1) + \exp(-i(\vec{k}_0 + \delta\vec{k})\vec{a}_2)| \\
&= \pm t \left| 1 + \exp\left(i\delta\vec{k}\frac{\sqrt{3}}{2}a\right) \cdot 2 \cos\left(\frac{2\pi}{3}a + \delta\vec{k}\frac{a}{2}\right) \right|
\end{aligned}$$

$$\approx \pm t \left| 1 - 1 - i \frac{\sqrt{3}}{2} a \delta \vec{k} \right| \approx \pm \frac{\sqrt{3}}{2} a t |\delta \vec{k}| \approx \frac{\sqrt{3}}{2} a t |\vec{k} - \vec{k}_0|$$

