

$$\vec{h} = \vec{F} = e(\vec{E} + \vec{v} \times \vec{B}) \quad \left. \begin{array}{l} \vec{v} = \frac{1}{\hbar} \frac{d\vec{z}(\vec{z})}{dt} \end{array} \right\} \text{kvaziklasični približek}$$

$$\textcircled{1} \vec{E} = 0, \vec{B}(\vec{r}) = (0, 0, B) \quad \rightarrow \quad \xi = \frac{\hbar^2 k^2}{2m}$$

proti elektroni:  $V(\vec{r}) = 0$  el. kroji s ciklotronske frekvence  $\frac{|e|B}{m} = \omega_c$

v periodičnem potencialu s lokalni min, max ali sedla

$$\xi(\vec{z}) = \frac{\hbar^2}{2} \vec{k} \cdot (m^*)^{-1} \vec{z} + \xi_0$$

$\hookrightarrow$  tenzor efektivne mase (simetričen)

$$\begin{aligned} \kappa_{\alpha} &= \frac{1}{\hbar} \frac{\partial \xi}{\partial k_{\alpha}} = \frac{1}{\hbar} \frac{\partial}{\partial k_{\alpha}} \left( \frac{\hbar^2}{2} k_{\beta} (m^*)^{-1}_{\beta\gamma} k_{\gamma} \right) = \frac{\hbar}{2} \left( \delta_{\alpha\beta} (m^*)^{-1}_{\beta\gamma} k_{\gamma} + k_{\beta} (m^*)^{-1}_{\alpha\gamma} \delta_{\beta\gamma} \right) = \\ &= \frac{\hbar}{2} \left( (m^*)^{-1}_{\alpha\gamma} k_{\gamma} + (m^*)^{-1}_{\alpha\beta} k_{\beta} \right) = \hbar (m^*)^{-1}_{\alpha\gamma} k_{\gamma} = \left( \hbar (m^*)^{-1} \vec{k} \right)_{\alpha} \end{aligned}$$

$$\vec{v} = \hbar (m^*)^{-1} \vec{k} \quad \rightarrow \quad m^* \vec{v} = \hbar \vec{k}$$

$$m^* \dot{\vec{v}} = e \vec{v} \times \vec{B}$$

$$m^* \dot{\vec{v}} = e \begin{pmatrix} v_x \\ v_y \\ v_z \end{pmatrix} \times \begin{pmatrix} 0 \\ 0 \\ B \end{pmatrix} = e \begin{pmatrix} v_y B \\ -v_x B \\ 0 \end{pmatrix}$$

$$m^* = \begin{pmatrix} m_{xx} & m_{xy} & m_{xz} \\ m_{yx} & m_{yy} & m_{yz} \\ m_{zx} & m_{zy} & m_{zz} \end{pmatrix}$$

$$\vec{v}(t) = \vec{v}_0 e^{-i\omega t} \quad \rightarrow \quad \vec{r}(t) = \int \vec{v}(t) dt$$

$$-i\omega m^* \vec{v}_0 = e \vec{v}_0 \times \vec{B}$$

$$m^* \vec{v}_0 = \frac{ie}{\omega} \begin{pmatrix} v_{0y} B \\ -v_{0x} B \\ 0 \end{pmatrix}$$

$$\begin{pmatrix} m_{xx} & m_{xy} - \frac{ieB}{\omega} & m_{xz} \\ m_{yx} + \frac{ieB}{\omega} & m_{yy} & m_{yz} \\ m_{zx} & m_{zy} & m_{zz} \end{pmatrix} \begin{pmatrix} v_{0x} \\ v_{0y} \\ v_{0z} \end{pmatrix} = 0$$

$$\det(\downarrow) = 0$$

$$m_{xx}m_{yy}m_{zz} + (m_{xy} - \frac{ieB}{\omega})m_{yz} + (m_{yx} + \frac{ieB}{\omega})m_{zy} - m_{xz}m_{yy}m_{zx} - m_{yz}m_{zy}m_{xx} - (m_{xy} - \frac{ieB}{\omega})(m_{yx} + \frac{ieB}{\omega})m_{zz} = 0$$

$$\det m^* + \left(\frac{ieB}{\omega}\right)^2 m_{zz} = 0$$

$$\det m^* = \frac{e^2 B^2}{\omega^2} m_{zz}$$

$$\omega^2 = \frac{e^2 B^2 m_{zz}}{\det m^*}$$

$$\omega = \frac{|e|B}{\sqrt{\frac{\det m^*}{m_{zz}}}}$$

$$m_{zz} = \hat{e} \cdot m^* \hat{e}$$

$$\omega = \frac{|e|B}{m_c^*}$$

$$m_c^* = \sqrt{\frac{\det m^*}{\hat{e} \cdot m^* \hat{e}}}$$

$\hookrightarrow$  ciklotronska efektivna masa  
 $m_c^* > 0$  s lokalni min. poun

$\hookrightarrow$  gostota stanj: proti elektroni:  $g(\xi) = \frac{1}{\pi^2 \hbar^3} \sqrt{2m^*} \sqrt{\xi} \approx 3D$

v periodičnem potencialu s lokalni min:

$$g(\xi) = \frac{1}{\pi^2 \hbar^3} \sqrt{2m_g^*} \sqrt{\xi - \xi_0}$$

$$m_g^* = 3 \sqrt{\det m^*}$$

Blochove oscilacije u 1D reznici  
u približni tesne rezi



$$\hbar \dot{k} = eE$$

$$k(t) = k(0) + \frac{eE}{\hbar} t$$

$$E > 0$$

$$E < 0$$

$$\epsilon(k) = -\gamma \sum_{R=n.s.f} e^{i\vec{z} \cdot \vec{k}} =$$

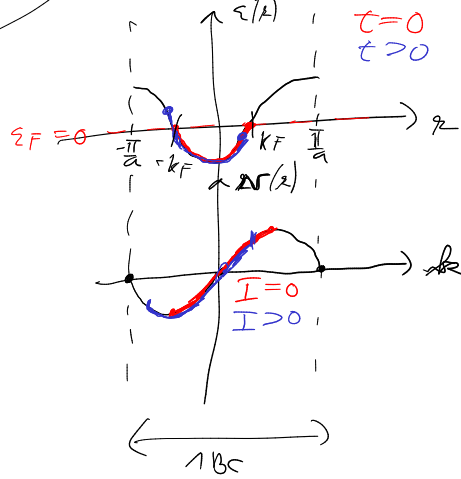
$$= -\gamma (e^{i2a} + e^{i2(-a)}) =$$

$$\boxed{\epsilon(k) = -2\gamma \cos 2a}$$

spin  
↓  
 $I(t) = e2 \sum_{|k(0)| < k_F} v(k(t)) \cdot \frac{1}{L}$  ( $j = mcv$ )

$$\Delta k = \frac{L}{2\pi}$$

$$v = \frac{1}{\hbar} \frac{d\epsilon}{dk} = \frac{2\gamma a \sin 2a}{\hbar}$$



$$I(t) = \frac{2e}{L} \sum_{|k(0)| < k_F} v(k(0) + \frac{eE}{\hbar} t) =$$

$$= \frac{2e}{L} \frac{L}{2\pi} \int_{-k_F}^{k_F} dk(0) v(k(0) + \frac{eE}{\hbar} t) =$$

$$= \frac{e}{\pi} \int_{-k_F}^{k_F} dk(0) \frac{2\gamma a}{\hbar} \sin(k(0) + \frac{eE}{\hbar} t)a =$$

$$= \frac{2e\gamma a}{\pi \hbar} \frac{1}{a} \left[ \cos(k(0) + \frac{eE}{\hbar} t)a \right]_{-k_F}^{k_F} =$$

$$= \frac{2e\gamma}{\pi \hbar} \left( \cos(k_F + \frac{eE}{\hbar} t)a - \cos(-k_F + \frac{eE}{\hbar} t)a \right) =$$

$$= -\frac{4e\gamma}{\pi \hbar} \sin 2F a \sin \frac{eE t a}{\hbar} =$$

$$= \frac{4e\gamma}{\pi \hbar} \sin 2F a \sin \frac{|eE| a}{\hbar} t =$$

$$= \frac{4e\gamma}{\pi \hbar} \sin 2F a \sin \omega_B t$$

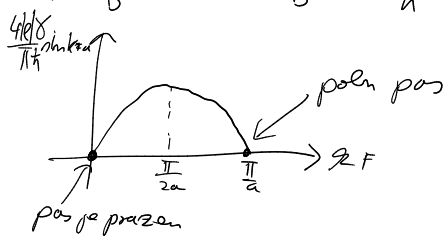
$$\omega_B = \frac{|eE| a}{\hbar}$$

Blochova frekvencija

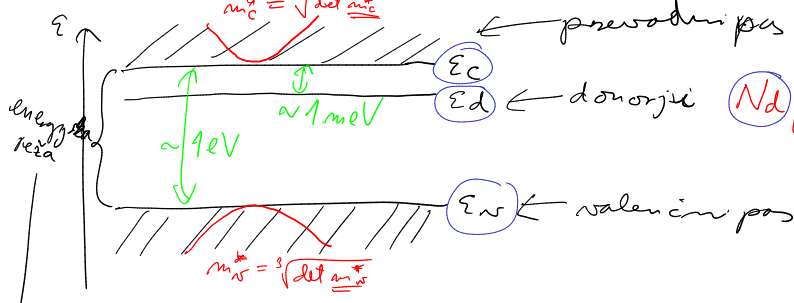
$$J \ll \frac{1}{\omega_B} \text{ (realne kvantne)}$$

čas med trkai

Ohmov zakon (Drude)



$\mu(T)$  v dopiranem polprevodniku



$$g_c(\epsilon) = \frac{\sqrt{2m_c^*}}{\pi^2 \hbar^3} \sqrt{\epsilon - \epsilon_c} \left[ \frac{1}{2m^3} \right]$$

$$g_v(\epsilon) = \frac{\sqrt{2m_v^*}}{\pi^2 \hbar^3} \sqrt{\epsilon_v - \epsilon} \left[ \frac{1}{2m^3} \right]$$

Konzentracija elektronov v prevodnem pasu

ohranitev nabojca:  $n_c = p_v + p_d \rightarrow$  vsajli v valeninem pasu vsajli na donornih nivojih  $\infty \rightarrow$  za dovolj velike T

$$n_c = \int_{\epsilon_c}^{\infty} d\epsilon g_c(\epsilon) f(\epsilon) = \int_{\epsilon_c}^{\infty} d\epsilon \frac{\sqrt{2m_c^*}}{\pi^2 \hbar^3} \sqrt{\epsilon - \epsilon_c} \frac{1}{e^{\beta(\epsilon - \mu)} + 1}$$

$\approx e^{-\beta(\epsilon - \mu)} \rightarrow$  mede generirani polprevodnik

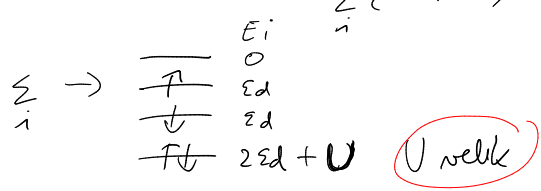
$$= \int_{\epsilon_c}^{\infty} d\epsilon \frac{\sqrt{2m_c^*}}{\pi^2 \hbar^3} \sqrt{\epsilon - \epsilon_c} e^{-\beta(\epsilon - \mu)} = \frac{\sqrt{2m_c^*}}{\pi^2 \hbar^3} \int_{\epsilon_c}^{\infty} d\epsilon \sqrt{\epsilon - \epsilon_c} e^{-\beta(\epsilon - \epsilon_c)} e^{-\beta(\epsilon_c - \mu)}$$

$$= \frac{\sqrt{2m_c^*}}{\pi^2 \hbar^3} \frac{\sqrt{\pi}}{2} \beta^{-3/2} e^{-\beta(\epsilon_c - \mu)} \rightarrow n_c = N_c e^{-\beta(\epsilon_c - \mu)}$$

$$N_c = \frac{1}{4} \left( \frac{2m_c^* k_B T}{\pi \hbar^2} \right)^{3/2}$$

$$p_v = \int_{-\infty}^{\epsilon_v} d\epsilon g_v(\epsilon) [1 - f(\epsilon)] = \dots = p_v = p_v e^{-\beta(\mu - \epsilon_v)} \quad p_v = \frac{1}{4} \left( \frac{2m_v^* k_B T}{\pi \hbar^2} \right)^{3/2}$$

1 donordni atom:  $\langle N \rangle = \frac{\sum_i N_i e^{-\beta(E_i - \mu N_i)}}{\sum_i e^{-\beta(E_i - \mu N_i)}} = \frac{0 e^{-\beta(0 - \mu)} + 2 \cdot 1 e^{-\beta(\epsilon_d - \mu)} + 2 e^{-\beta(2\epsilon_d + U - \mu)}}{e^{-\beta(0 - \mu)} + 2 e^{-\beta(\epsilon_d - \mu)} + e^{-\beta(2\epsilon_d + U - \mu)}}$



zanemarimo

$$\langle N \rangle = \frac{2e^{-\beta(\epsilon_d - \mu)}}{1 + 2e^{-\beta(\epsilon_d - \mu)}}$$

$$p_d = N_d (1 - \langle N \rangle)$$

$$p_d = N_d \frac{1}{1 + 2e^{-\beta(\epsilon_d - \mu)}}$$

$n_c = p_v + p_d$

$$N_c e^{-\beta(\epsilon_c - \mu)} = p_v e^{-\beta(\mu - \epsilon_v)} + \frac{N_d}{1 + 2e^{-\beta(\epsilon_d - \mu)}} \rightarrow$$

polinom 3 stopnje za nemarimo  $e^{\beta\mu}$ .

zanemarimo pri nizki temperaturah  $\rightarrow$  kvadratna enačba za  $e^{\beta\mu}$

$$N_c e^{-\beta\epsilon_c} e^{\beta\mu} + 2N_c e^{-\beta(\epsilon_c + \epsilon_d)} e^{2\beta\mu} - N_d = 0$$

$$e^{\beta\mu} = \frac{-N_c e^{-\beta\varepsilon_c} \pm \sqrt{(N_c e^{-\beta\varepsilon_c})^2 + 8N_c N_d e^{-\beta(\varepsilon_c + \varepsilon_d)}}}{4N_c e^{-\beta(\varepsilon_c + \varepsilon_d)}}$$

$$e^{\beta\mu} = -\frac{1}{4} e^{\beta\varepsilon_d} \pm \sqrt{\left(\frac{1}{4} e^{\beta\varepsilon_d}\right)^2 + \frac{N_d}{2N_c} e^{\beta(\varepsilon_c + \varepsilon_d)}}$$

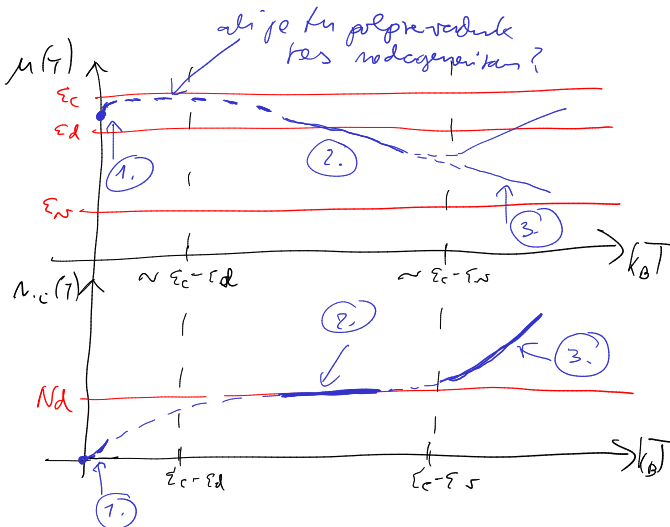
$$e^{\beta\mu} = \frac{1}{4} e^{\beta\varepsilon_d} \left[ -1 \oplus \sqrt{1 + \frac{8N_d}{N_c} e^{\beta(\varepsilon_c - \varepsilon_d)}} \right] \quad \beta = \frac{1}{k_B T}$$

1.  $\frac{8N_d}{N_c} e^{\beta(\varepsilon_c - \varepsilon_d)} \gg 1$     nizka temperatura     $k_B T \ll A(\varepsilon_c - \varepsilon_d)$      $A = \sigma(\eta)$      $N_c \propto T^{3/2}$ ;  $\varepsilon_c - \varepsilon_d > 0$

$$e^{\beta\mu} = \frac{1}{4} e^{\beta\varepsilon_d} \sqrt{\frac{8N_d}{N_c}} e^{\beta \frac{\varepsilon_c - \varepsilon_d}{2}} = \sqrt{\frac{N_d}{2N_c}} e^{\beta \frac{\varepsilon_c + \varepsilon_d}{2}} \quad \lim_{T \rightarrow 0} = 0$$

$$\mu = \frac{\varepsilon_c + \varepsilon_d}{2} + \frac{1}{\beta} \ln \sqrt{\frac{N_d}{2N_c}} = \frac{\varepsilon_c + \varepsilon_d}{2} - \frac{3}{4} k_B T \ln T + \sigma(T)$$

$$n_c = N_c e^{-\beta(\varepsilon_c - \mu)} = N_c e^{-\beta\varepsilon_c} \cdot \sqrt{\frac{N_d}{2N_c}} e^{\beta \frac{\varepsilon_c + \varepsilon_d}{2}} = \sqrt{\frac{N_d N_c}{2}} e^{-\beta \frac{\varepsilon_c - \varepsilon_d}{2}} \propto T^{3/4} e^{-\frac{\varepsilon_c - \varepsilon_d}{2k_B T}}$$



2.  $\frac{8N_d}{N_c} e^{\beta(\varepsilon_c - \varepsilon_d)} \ll 1$     visoka temperatura     $k_B T \gg A(\varepsilon_c - \varepsilon_d)$      $A = \sigma(\eta)$

$$e^{\beta\mu} = \frac{1}{4} e^{\beta\varepsilon_d} \left[ -1 + \left( 1 + \frac{1}{2} \frac{8N_d}{N_c} e^{\beta(\varepsilon_c - \varepsilon_d)} \right) \right] = \frac{N_d}{N_c} e^{\beta\varepsilon_c}$$

$$\mu = \varepsilon_c + k_B T \ln \frac{N_d}{N_c} = \varepsilon_c - \frac{3}{2} k_B T \ln T + \sigma(T)$$

$$n_c = N_c e^{-\beta\varepsilon_c} e^{\beta\mu} = N_c e^{-\beta\varepsilon_c} \frac{N_d}{N_c} e^{\beta\varepsilon_c} = N_d$$

3.  $n_c \gg N_d$     ( $T \gtrsim \varepsilon_c - \varepsilon_s$ )

$$n_c = p_n + p_d$$

$$N_c e^{-\beta(\varepsilon_c - \mu)} = p_n e^{-\beta(\mu - \varepsilon_s)}$$

$$e^{2\beta\mu} = \frac{p_n}{N_c} e^{\beta(\varepsilon_c + \varepsilon_s)}$$

$$N_c = \frac{1}{4} \left( \frac{2 m_c^* k_B T}{\pi \hbar^2} \right)^{3/2}$$

$$p_n = \frac{1}{4} \left( \frac{2 m_n^* k_B T}{\pi \hbar^2} \right)^{3/2}$$

$$\mu = \frac{\varepsilon_c + \varepsilon_s}{2} + \frac{1}{2} k_B T \ln \frac{p_n}{N_c} = \frac{\varepsilon_c + \varepsilon_s}{2} + \frac{3}{4} k_B T \ln \frac{m_n^*}{m_c^*}$$

$$\begin{aligned}
 m_c &= N_c \overset{\downarrow}{e^{-\beta \epsilon_c}} e^{\beta \mu} = N_c e^{-\beta \epsilon_c} \sqrt{\frac{P_N}{N_c}} e^{\beta \frac{\epsilon_c + \epsilon_N}{2}} = \\
 &= \sqrt{N_c P_N} e^{-\beta \left( \frac{\epsilon_c - \epsilon_N}{2} \right)} \alpha T^{3/2} e^{-\frac{\epsilon_c - \epsilon_N}{2 k_B T}}
 \end{aligned}$$