

Elektronika

①

- elektron (grška beseda za jantar). Iz drobnostim
kristala jantarja z blagom in jantar naelektrici. 600 BC
Jantar lahko privlači delčke nekega črna-slame.
- Tedaj so poznali tudi želez magnet (Magnet) grška
pohranina. Vendar niso poznali povezane med magnetom
ter el tokom tri do leta 1820
- 1820 Hans Christian Oersted - tok po prevodniku
lahko ukloni magnetno iglo

Naloga

- Poznamo dve vrsti naloge: pozitivni in negativni
- Prikladi med rozvostrnimi ter odlog med istostrnimi
nalogi
- Ohranitev naloge $(l_1) (l_2) \begin{array}{c} | \\ \text{el} \\ \text{nis} \\ \text{oo} \end{array} \Rightarrow l_1 + l_2 = l$

Prevodniki

Isolatorji

Stevo - del elektronov je
skrajno prava gibljivi

Elektroni niso prava
gibljivi - vezani na pozitivne
ione

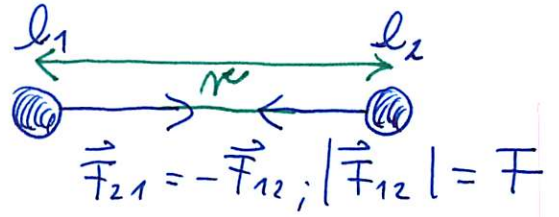
- Naloga je kvantitativna: $l_0 = 1,6 \cdot 10^{-19} \text{ As}$ Amper!
André-Marie
(1775-1836)

Coulombov zakon

(2)

Električna sila med dvema točkastima naboja:

$$F = \frac{q_1 q_2}{4\pi \epsilon_0 r^2}$$



ϵ_0 : dielektrična konstanta = $8,9 \cdot 10^{-12} \text{ As/Vm}$

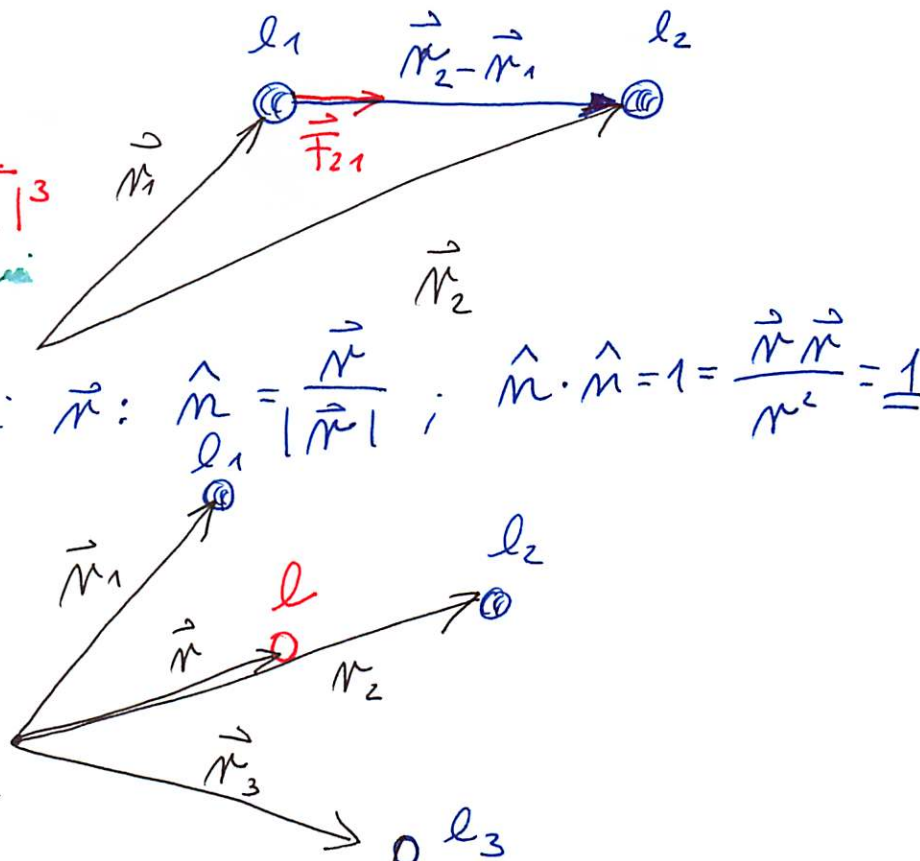
q_1, q_2 električna naboja v enotah [As]

Zapis z vektorji:

$$\vec{F}_{21} = \frac{q_1 q_2}{4\pi \epsilon_0} \frac{\vec{r}_1 - \vec{r}_2}{|\vec{r}_1 - \vec{r}_2|^3}$$

Sila drugega naboja na prvi naboju.

Enotski vektor v smeri \vec{r} :



$$\hat{r} = \frac{\vec{r}}{|\vec{r}|}; \quad \hat{r} \cdot \hat{r} = 1 = \frac{\vec{r} \cdot \vec{r}}{r^2} = \underline{1}$$

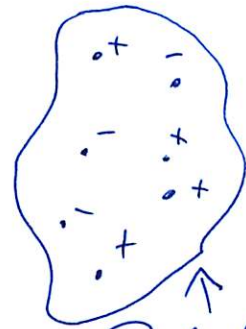
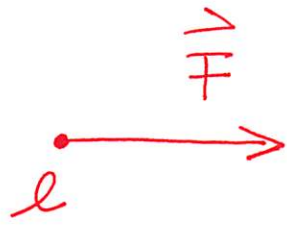
$$\vec{F}_{\vec{r}} = \sum_i \frac{q q_i}{4\pi \epsilon_0} \frac{\vec{r} - \vec{r}_i}{|\vec{r} - \vec{r}_i|^3}$$

teči po vseh nabojih q_1, q_2, \dots ,
oznan po \vec{r} !

Sila na izbrani naboju q v točki \vec{r} kot posledica vseh ostalih nabojev.

Jakost električnega polja

(3)



Poljubna porazdelitev nabojev

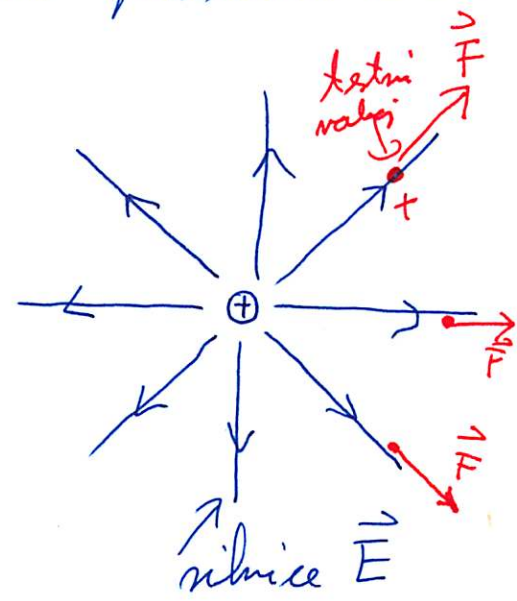
Električno polje povzroči silo na naboju:

$$\vec{F} = q \vec{E}$$

E [$\frac{V}{m}$] enota: V - Volt (Alessandro Volta 1745-1827)
Italijanski fizik

Izumitelj prve kemične baterije

Električno polje obdaja nabite delce. Smer silnice \vec{E} polja napoveduje smer sile na testni pozitivni električni naboju.



Električno polje točkastega naboja:

$$E = \frac{q}{4\pi\epsilon_0 r^2}$$

vektor:

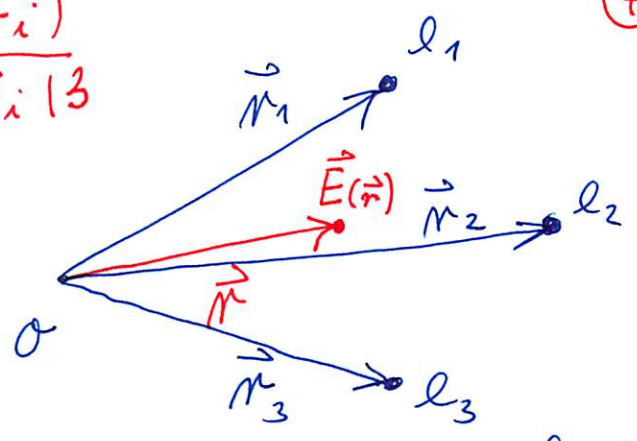
$$\vec{E}(\vec{r}) = \frac{q}{4\pi\epsilon_0 r^2} \vec{e}_r$$

če izhodišče koordinatnega sistema v $\vec{r} = 0$! meri vektora v radialni smeri.

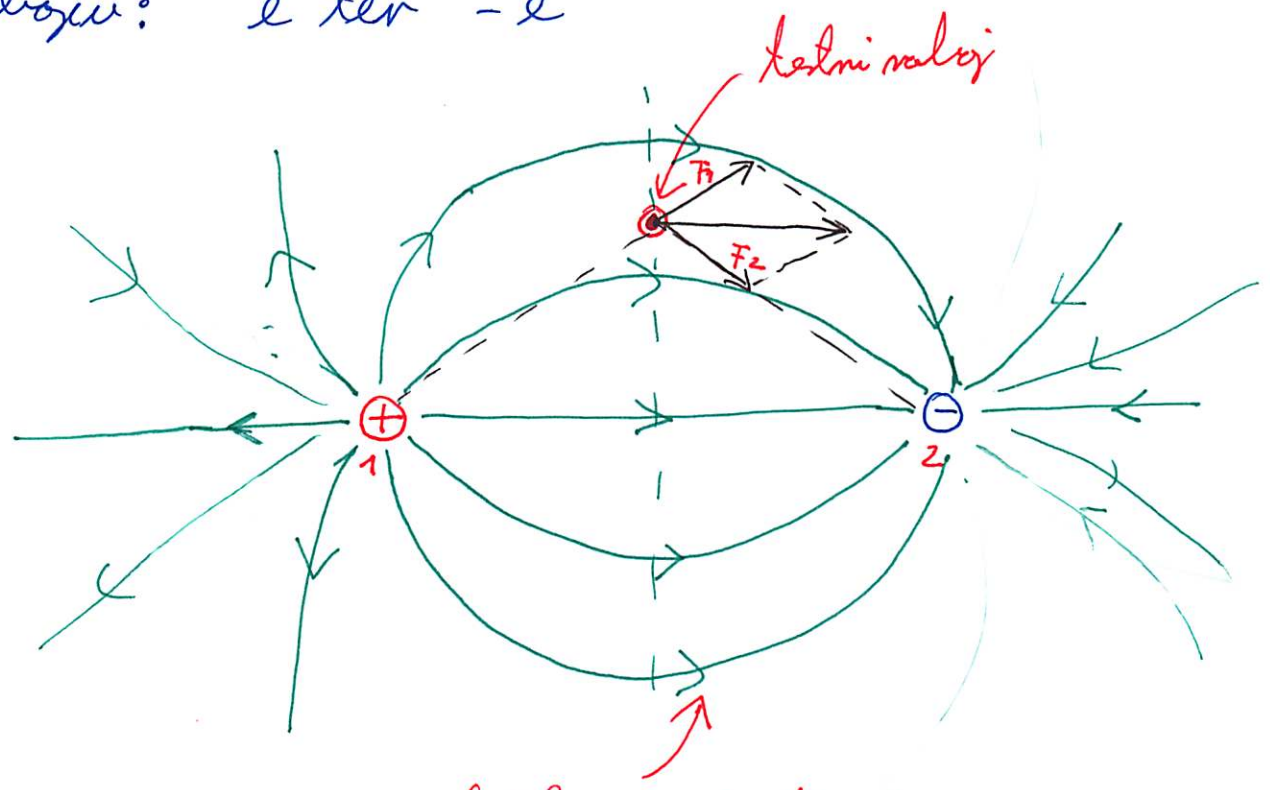
Primeri \vec{E} polj različnih točkastih nabojev se sestavajo!

$$\vec{E}(\vec{r}) = \sum_i \frac{q_i}{4\pi\epsilon_0} \frac{(\vec{r} - \vec{r}_i)}{|\vec{r} - \vec{r}_i|^3}$$

i
po vsak
naboju

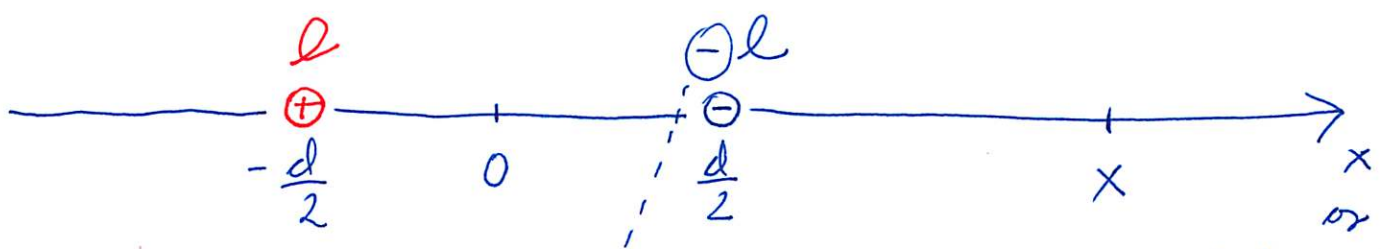


Slika nabitih v okolici dveh nasprotno enakih naboju: e ter $-e$



Električni dipol

El. polje elekt. dipola



$$E(x) = \frac{l}{4\pi\epsilon_0} \left(\frac{1}{(x + \frac{d}{2})^2} - \frac{1}{(x - \frac{d}{2})^2} \right) = \frac{l}{4\pi\epsilon_0} \frac{-2dx}{(x^2 - \frac{d^2}{4})^2} \Rightarrow$$

$$E(x) \approx \frac{-l}{4\pi\epsilon_0} \frac{2dx}{x^4 - \frac{x^2 d^2}{2} + \dots}$$

Zanemarimo da $x \gg d!$

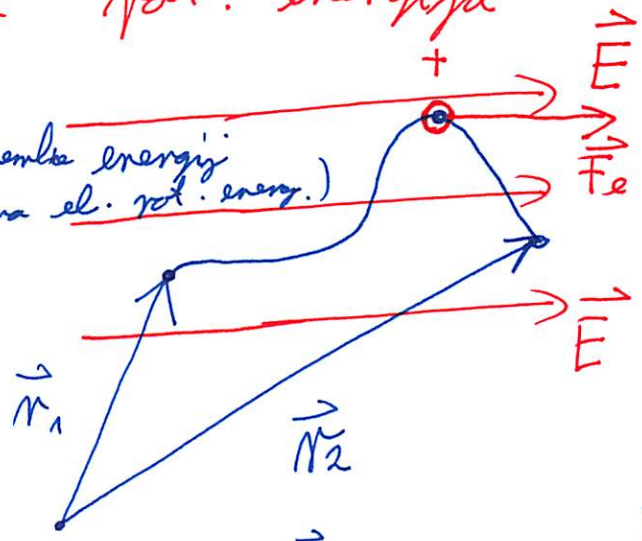
$$E(x) = - \frac{ed}{2\pi\epsilon_0 x^3} = - \frac{\mu}{2\pi\epsilon_0 x^3}$$

μ : elektrinski dipolni moment!

Elektrinska pot. energija

$A + A_{el} = \Delta W$
 \uparrow delo elektrinske sile
 delo vseh ostalih sil

← sprememba energij
 (na zojima el. pot. energ.)



$$A = - \underbrace{A_{el}}_{\Delta W_{el}} + \Delta W$$

$$A_{el} = \int_{\vec{r}_1}^{\vec{r}_2} \vec{F}_e d\vec{s} = e \int_{\vec{r}_1}^{\vec{r}_2} \vec{E} d\vec{s}$$

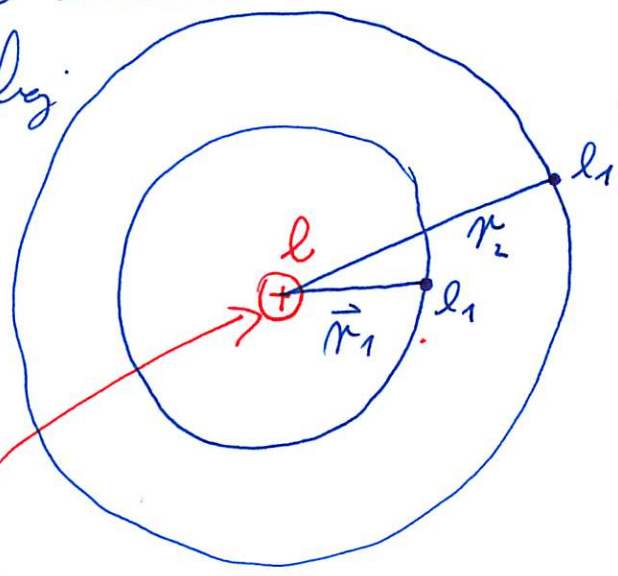
$$\Delta W_{el} = - e \int_{\vec{r}_1}^{\vec{r}_2} \vec{E} d\vec{s} = e \underbrace{U(\vec{r}_2, \vec{r}_1)}_{\text{napetost med } \vec{r}_2 \text{ in } \vec{r}_1} = \underbrace{eV(\vec{r}_2) - eV(\vec{r}_1)}_{\text{elektrinska potencialna energija}}$$

$W_{el} = q V$; $V(\vec{r}) \doteq$ elektrini potencial

$U(\vec{r}_2, \vec{r}_1) = V(\vec{r}_2) - V(\vec{r}_1)$

Napetost je vedro definirana med dvema točkama! Enote: U in V [V = Volt]

Primer: Točkasti nabo



Sprememba el. pot. energije naloga q_1 v polju naloga

q_1 :

$$\Delta W_{el} = -q_1 \int_{r_1}^{r_2} \frac{q}{4\pi \epsilon_0 r^2} dr = q_1 \left[\frac{q}{4\pi \epsilon_0 r_2} - \frac{q}{4\pi \epsilon_0 r_1} \right]$$

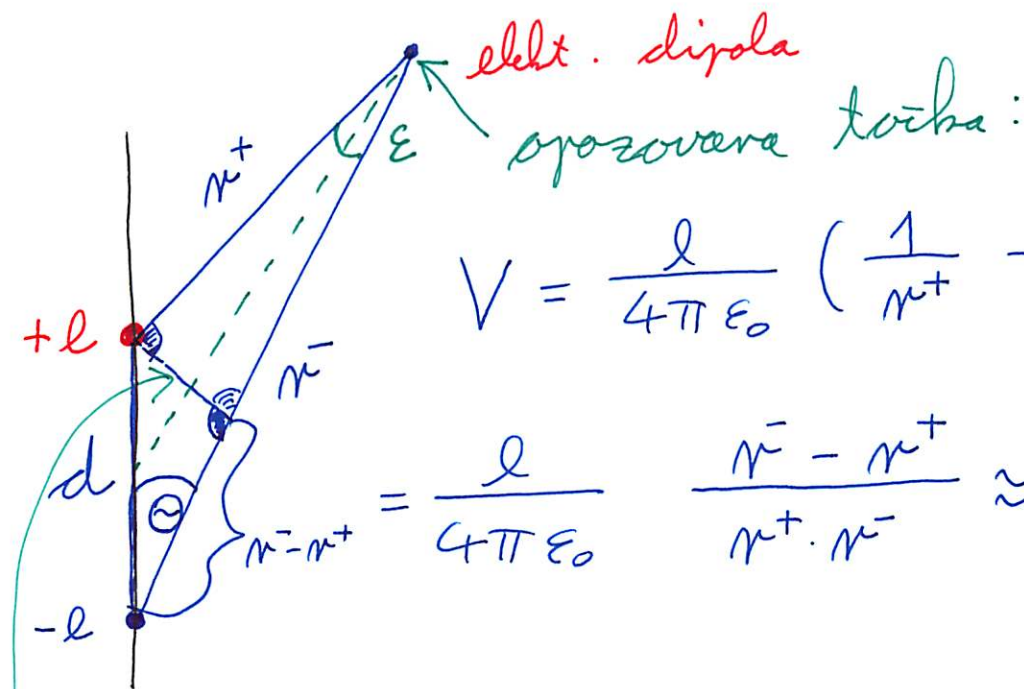
$= q_1 (V(r_2) - V(r_1))$, pri čemer velja:

$$V(r) = \frac{q}{4\pi \epsilon_0 r}$$

Elektrini potencial elektrinsga točkastega naboja!

Električni potencial

(7)



$$V = \frac{l}{4\pi\epsilon_0} \left(\frac{1}{r^+} - \frac{1}{r^-} \right) =$$

$$= \frac{l}{4\pi\epsilon_0} \frac{r^- - r^+}{r^+ \cdot r^-} \approx \frac{l d \cos\theta}{4\pi\epsilon_0 r^2}$$

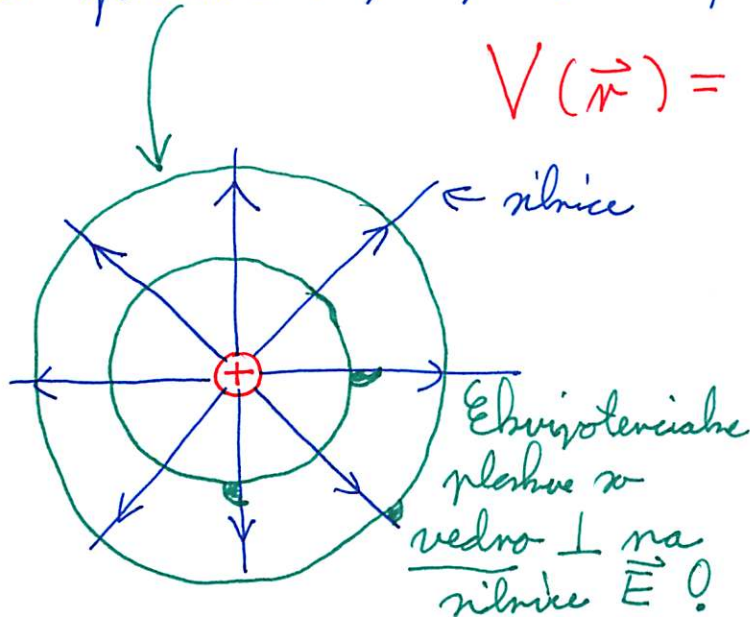
"prava točka" le približno, saj gre za evakobinski
trikotnik pri kotezi $d \ll r^+, r^-$ oziroma $\epsilon \ll \frac{\pi}{2}$
↓ električni dipol

$$V(r, \theta) = \frac{p \cos\theta}{4\pi\epsilon_0 r^2}$$

Elektripotencialne ploskve:

So ploskve kjer je el. potencial konstanten:

$$V(\vec{r}) = V_0 \text{ (konstanta)}$$



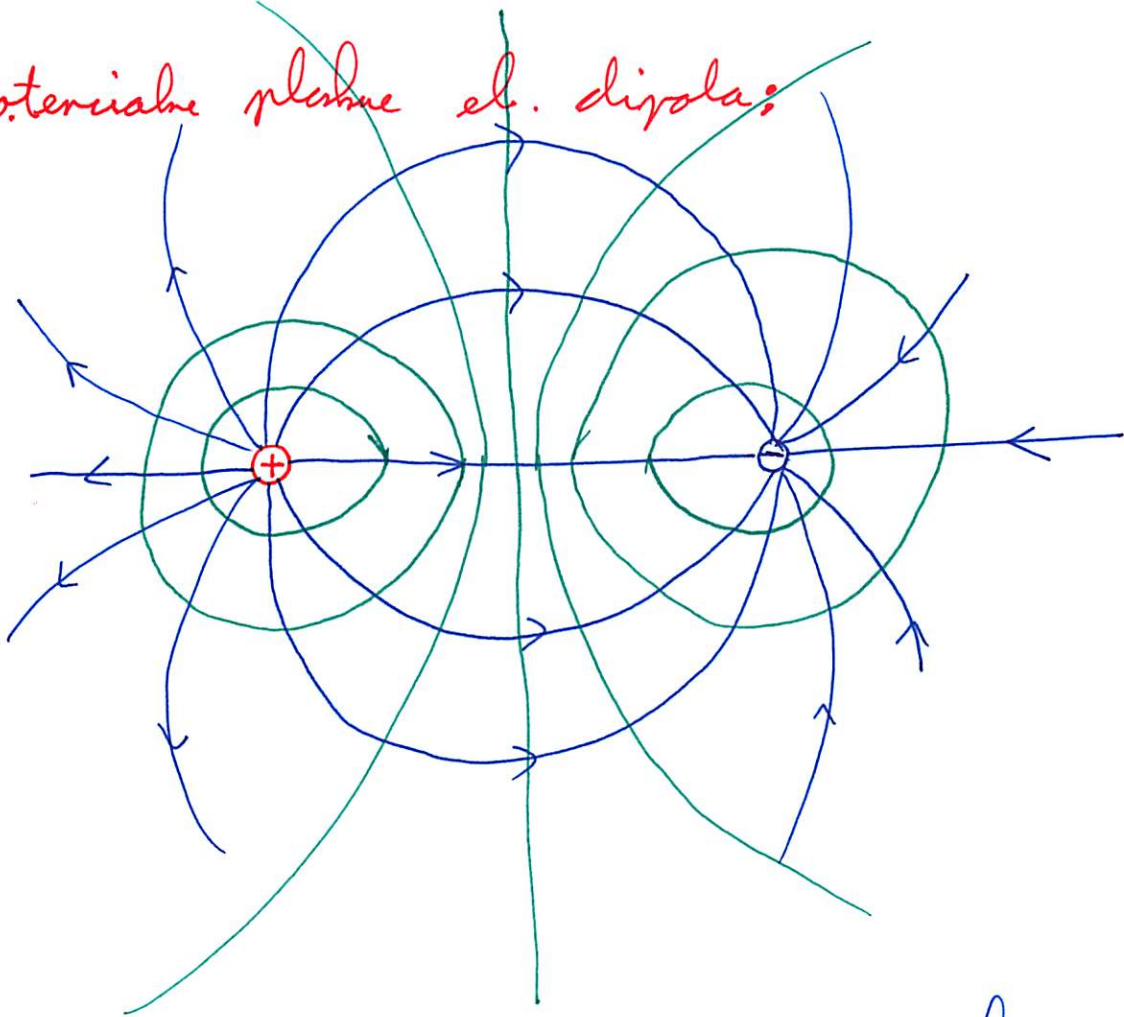
$$\frac{l}{4\pi\epsilon_0 r} = V_0 = \text{konstanta}$$

velja če je $r = r_0$!

Ta pa je kroglja z radijem r_0 !

Ekvipotencijalne ploskve el. dipola:

8



Če se nahaja delček giblje vzdolž ekvipotencialnih ploskev, se mu ne spremeni električna potencialna energija.

Rešitevno: Če dajmo izračunamo \vec{E} , če poznamo $V(\vec{r})$?

$$-\int_{\vec{r}_1}^{\vec{r}_2} \vec{E} d\vec{r} = V(\vec{r}_2) - V(\vec{r}_1) \Rightarrow \vec{E} = - \underset{\substack{\uparrow \\ \text{gradient}}}{\nabla} V(\vec{r})$$

$$\vec{\nabla} = \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \text{ oziroma } \vec{\nabla} V(r) = \left(\frac{\partial V}{\partial x}, \frac{\partial V}{\partial y}, \frac{\partial V}{\partial z} \right)$$

$$\text{Primer: } V(r) = \frac{q}{4\pi\epsilon_0 r} ; \frac{\partial V}{\partial x} = \frac{\partial V}{\partial r} \frac{\partial r}{\partial x}$$

$$r = \sqrt{x^2 + y^2 + z^2} ; \frac{\partial r}{\partial x} = \frac{2x}{2\sqrt{x^2 + y^2 + z^2}} = \frac{x}{r}$$

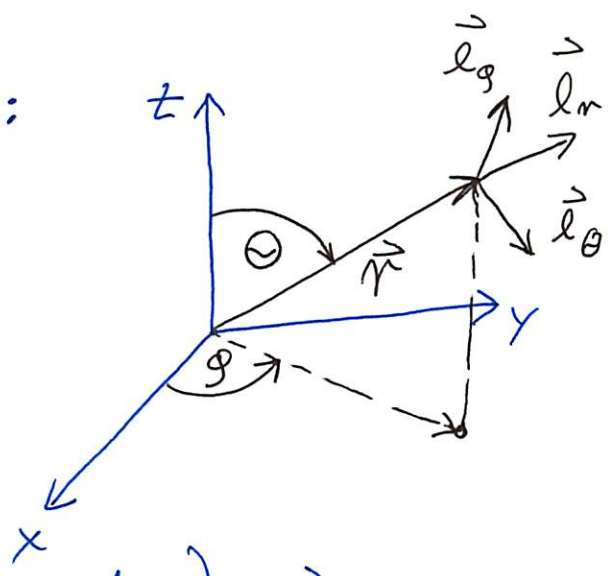
$$\vec{E} = - \vec{\nabla} V(r) = - \frac{q}{4\pi\epsilon_0} \underbrace{\frac{\partial}{\partial r}}_{-\frac{1}{r^2}} \left(\frac{\partial r}{\partial x}, \frac{\partial r}{\partial y}, \frac{\partial r}{\partial z} \right) =$$

$$= \frac{q}{4\pi\epsilon_0 r^2} \left(\frac{x}{r}, \frac{y}{r}, \frac{z}{r} \right) \Rightarrow$$

$$\boxed{E = \frac{q}{4\pi\epsilon_0 r^2} \frac{\vec{r}}{r}} \quad \checkmark$$

Drugi primer:

$$V(\vec{r}) = \frac{q d \cos \theta}{4\pi\epsilon_0 r^2};$$



$$\vec{\nabla} = \frac{\partial}{\partial r} \vec{l}_r + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} \vec{l}_\phi + \frac{1}{r} \frac{\partial}{\partial \theta} \vec{l}_\theta$$

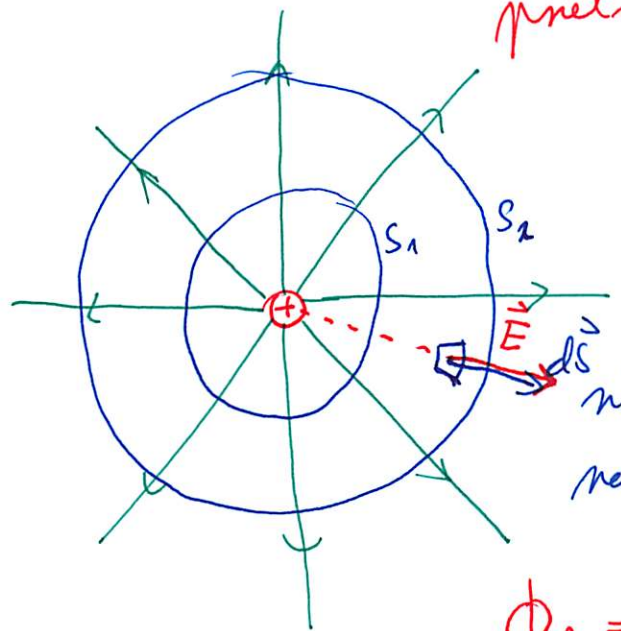
ali

$$\vec{\nabla} V = \frac{\partial V}{\partial r} \vec{l}_r + \frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi} \vec{l}_\phi + \frac{1}{r} \frac{\partial V}{\partial \theta} \vec{l}_\theta; \text{ ta } V(r^2)$$

$$= - \frac{q}{2\pi\epsilon_0} \frac{\cos \theta}{r^3} \vec{l}_r - \frac{q}{2\pi\epsilon_0 r^3} \sin \theta \vec{l}_\theta$$

$$\vec{E}(r, \theta) = \frac{q}{2\pi\epsilon_0 r^3} \left(\cos \theta \vec{l}_r + \sin \theta \vec{l}_\theta \right)$$

Zakon o električnem pretoku

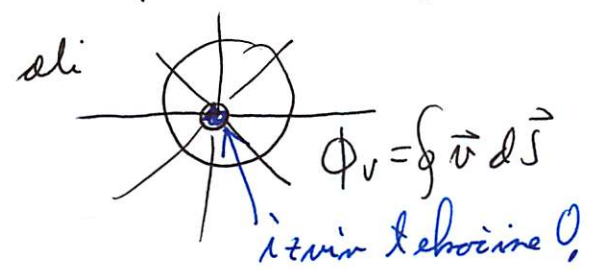


Število silnic, ki prebode S_1 ali S_2 je enako.
 Ali iz oblike silnic poročelne silnice lahko sklepamo na poročelne površine?

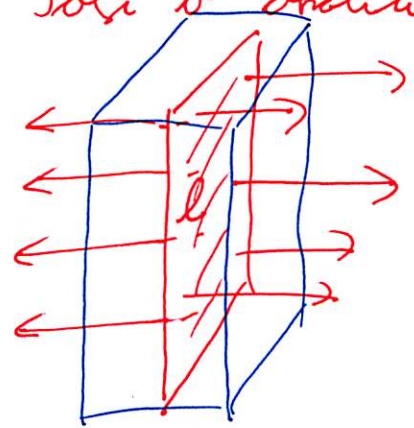
$$\Phi_e = \epsilon_0 \oint \vec{E} \cdot d\vec{S} = q$$

električni pretok je enak nalozi, ki ga objamimo z razširjenim ploskvinjo

Podobno kot tokovnice pri pretoku tekočine, si lahko mislimo, da tudi el. silnice predstavljajo tok električnega polja.
 $\Phi_v = \vec{v} \cdot \vec{S}$; $\Phi_e = \epsilon_0 \oint \vec{E} \cdot d\vec{S}$



Polji v obliki ravne plosvine:

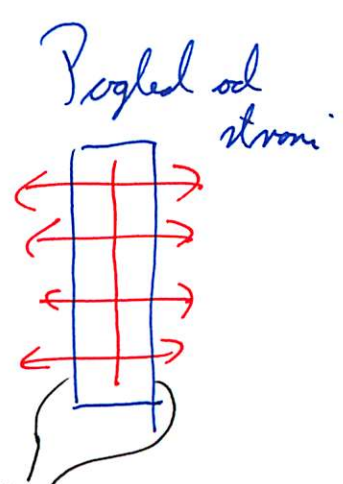


$$q = \epsilon_0 \oint \vec{E} \cdot d\vec{S}$$

$$q = \epsilon_0 E \cdot 2S$$

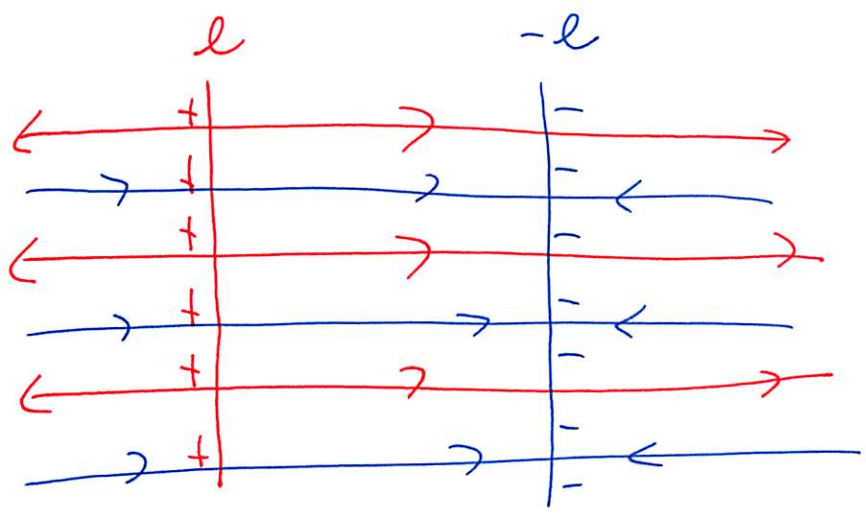
$$E = \frac{q}{\epsilon_0 2S} = \frac{\sigma}{2\epsilon_0}$$

$\sigma = \frac{q}{S}$: površinska gostota naložbe.



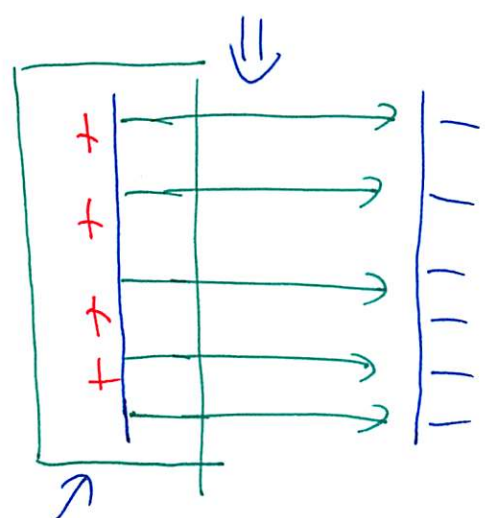
due strani $\Rightarrow 2S$

Ylandersaton



$$E = \frac{\sigma}{2\epsilon_0} - \frac{-\sigma}{2\epsilon_0} = \frac{\sigma}{\epsilon_0}$$

Primerak (+) plati
Primerak (-) plati

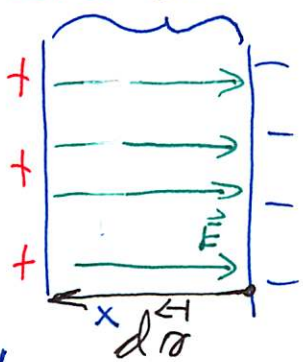


rezultati: $E = \frac{l}{\epsilon_0 S}$

Zunji ni \vec{E} polja!

alternativna izvedba: $l = \epsilon_0 \oint \vec{E} d\vec{S} = \epsilon_0 E S$
 $\Rightarrow E = \frac{l}{\epsilon_0 S}$ le en S, raj
 Zunji ni el. polja!

Kajetost med plati d
 kondensatorja?



$$\Rightarrow U = Ed = \frac{l d}{\epsilon_0 S}$$

$$U = - \int_{\ominus}^{\oplus} \vec{E} d\vec{r} = Ed$$

druga plati

$$U = - \int_{\ominus}^{\oplus} \vec{E} d\vec{r} = \int_0^d E dx = Ed$$

$$l = C U; \quad C = \frac{\epsilon_0 S}{d}$$

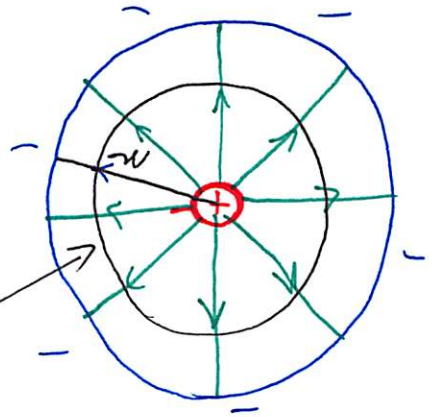
Ypociteta
 kondensatorja

Elektrische Felder in Kondensatoren

$$Q = \epsilon_0 \int \vec{E} \cdot d\vec{S} = \epsilon_0 E 2\pi r L$$

$E(r)$ ist konstant im Inneren r

L : Länge Kondensators



$$E = \frac{2\pi \epsilon_0 n Q}{L}$$

Weg zu E geht
 in Problem
 Kondensator?

$$Q = - \int_n^R \vec{E} \cdot d\vec{n} = - \frac{Q}{L} \int_n^R \frac{1}{r} dr = - \frac{Q}{L} \ln \frac{R}{n}$$

$$Q = \frac{2\pi \epsilon_0 L}{\ln \frac{R}{n}} Q \Rightarrow C = \frac{2\pi \epsilon_0 L}{\ln \frac{R}{n}}$$

Elektrische Felder
 Kondensator

Abbau Kondensator

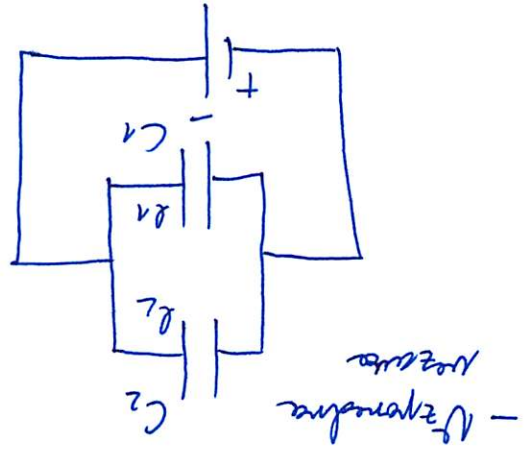
$$C_n = C_1 + C_2$$

Wachstum Kapazität

$$R_1 = C_1 U$$

$$R_2 = C_2 U$$

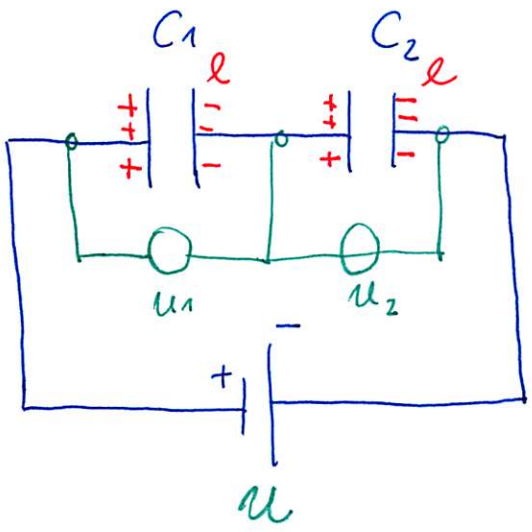
$$R_1 + R_2 = (C_1 + C_2) U = Q$$



- Zline, hat da kein
 vertikal gerinne!

relektm
 noly, na
 gleich Kondensat.

Zajedno vezava:



$$Q = C_1 U_1$$

$$Q = C_2 U_2$$

$$\frac{Q}{C_1} + \frac{Q}{C_2} = U_1 + U_2 = U$$

$$\frac{Q}{C_m} = U \Rightarrow$$

Kapacitet se razdeli: $U = U_1 + U_2$

$$C_m = \left(\frac{1}{C_1} + \frac{1}{C_2} \right)^{-1}$$

Energija kondenzatorja:

$A = e U$; delo je enako spremeni el. pot. energiji

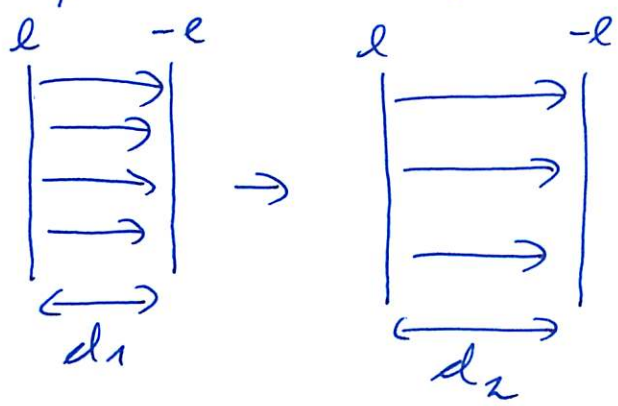
Toda $U = U(e)$; saj velja $U = \frac{e}{C}$

$$\Rightarrow A = \int_{e_1}^{e_2} U de = \int_{e_1}^{e_2} \frac{e de}{C} = \frac{e_2^2}{2C} - \frac{e_1^2}{2C} = W_{C_2} - W_{C_1}$$

$$W_C = \frac{e^2}{2C} = \frac{C}{2} U^2 = \frac{Ue}{2}$$

Energija kondenzatorja (nabitega)

Primer: bolika dela opravimo, a razmaknimo dve izolirani, nabiti plati kondenzatorja

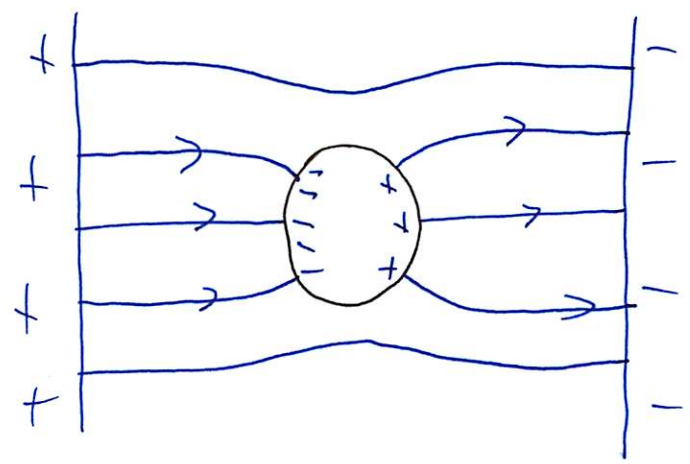


$$\begin{aligned}
 A &= W_{C_2} - W_{C_1} = \\
 &= \frac{e^2}{2} \left(\frac{1}{C_2} - \frac{1}{C_1} \right) = \\
 &= \frac{e^2}{2 \epsilon_0 S} (d_2 - d_1)
 \end{aligned}$$

Snov v električnem polju:

a.) Prevodnik

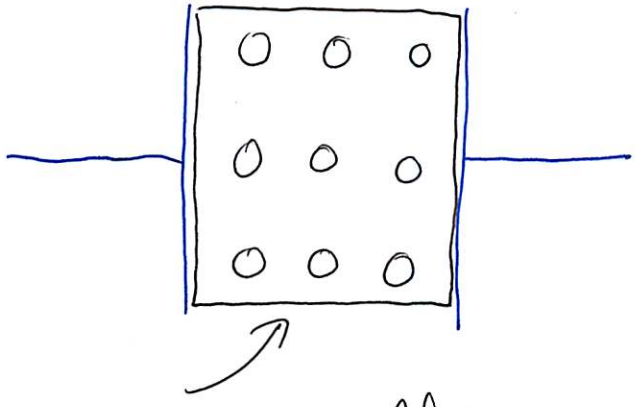
- V bravi ni \vec{E} polja in ni el. toka
- Silnice so \perp na površini brave



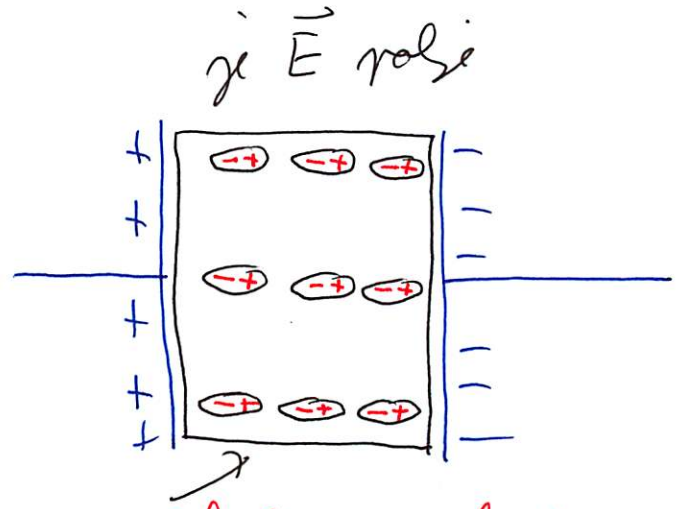
- V notranjosti brave ni \vec{E} polja (zovadi influence) in in ni naloživ!

b.) Izolator:

ni \vec{E} polja



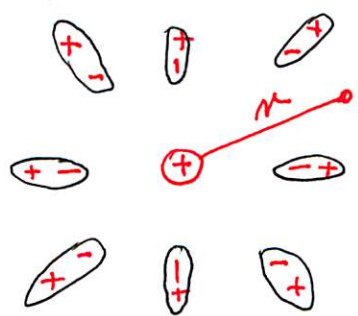
nepolarne molekule



si \vec{E} polje

el. nalož v dielektriku

molekule se polarizirajo \Rightarrow
 \vec{E} polje se "zorenči" ali



$$\vec{E} = \frac{q}{4\pi\epsilon\epsilon_0 N^2}$$

Zmanjša

$$\vec{E} = \frac{q}{\epsilon\epsilon_0 S}$$

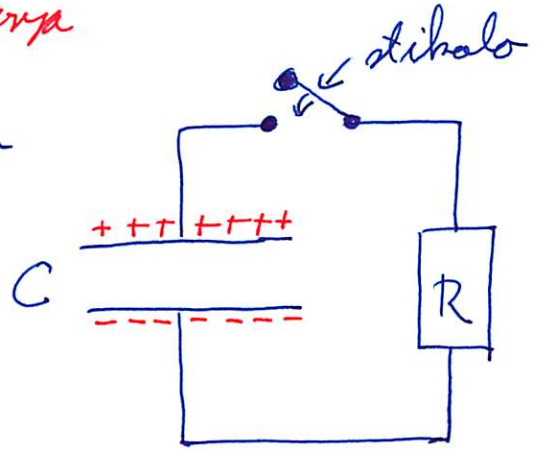
ϵ : dielektricitet

snov	ϵ
steklo	5-7
voda	81

Proznenji in polnenji kondenzatorja

- Proznenji

Uklo hitro se sprozi kondenzator preko upora in določi $I(t)$

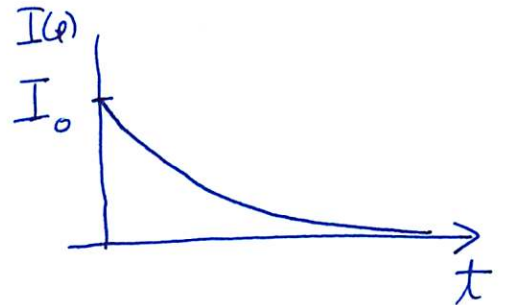


$$U_C + U_R = 0$$

$$\frac{q}{C} + IR = 0 \quad | \frac{d}{dt} I$$

$$\frac{I}{C} + \frac{dI}{dt} R = 0 \Rightarrow \int_{I_0}^I \frac{dI}{I} = - \int_0^t \frac{dt}{RC} = - \frac{t}{RC}$$

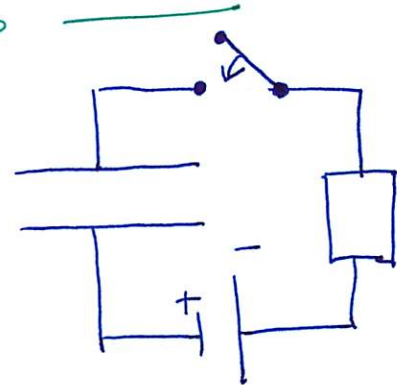
$$\ln \frac{I}{I_0} = - \frac{t}{RC} \Rightarrow I = I_0 e^{-\frac{t}{RC}}$$



Udeležba I_0 ?

$$I(t=0) = I_0 \Rightarrow \boxed{I_0 = -\frac{U}{RC}}$$

- Polnenje:



$$U = IR + \frac{q}{C}; \quad U = konst$$

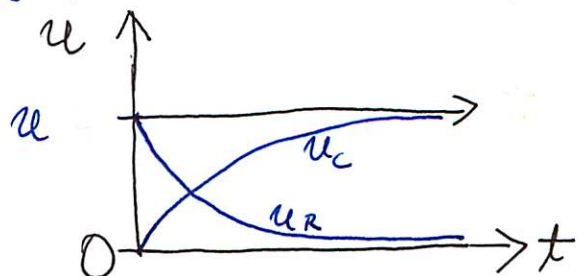
$$0 = \frac{dI}{dt} R + \frac{I}{C} \Rightarrow I = I_0 e^{-\frac{t}{RC}}$$

$$U_C = U - IR = U - I_0 R e^{-\frac{t}{RC}}$$

$$U_C(t=0) = 0 \Rightarrow U_C(t=0) = U - I_0 R = 0 \Rightarrow I_0 = \frac{U}{R}$$

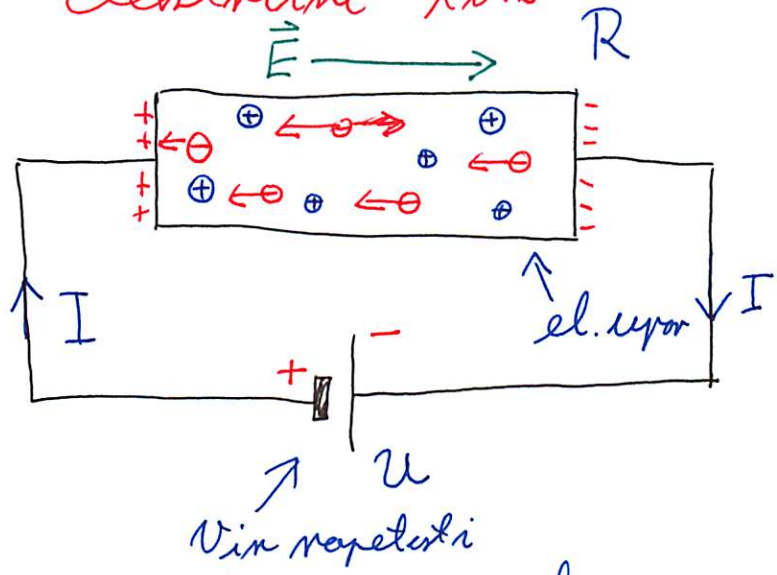
$$U_C = U (1 - e^{-\frac{t}{RC}})$$

$$U_R = U e^{-\frac{t}{RC}}$$



Električni tok

Model



- Po vezju teče električni tok I

$$I = \frac{dq}{dt} \left[\frac{A \cdot r}{r} = A \right]; \quad I: \text{množična veličina, ki gre}$$

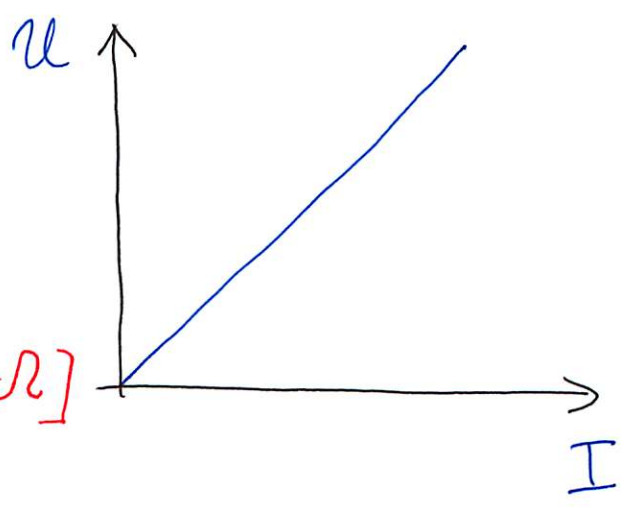
skozi preseč žice (upor) ne ta sama enota

- V snovi (prevodniku) po katerem teče el. tok je električno polje \vec{E}

- V točki snovi se premikajo le elektroni z negativnim nabojem tem v obratni smeri električnega toka

- Dejstvo konstantni nile va \vec{e}^- se \vec{e}^- v povprečju premikajo enakomerno!

$$U = I R$$



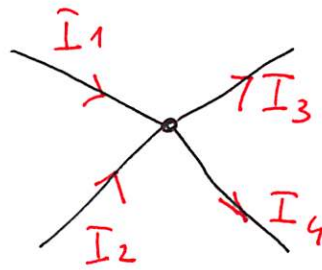
Ohmova zakon; $R \left[\frac{V}{A} = \Omega \right]$

Ω : 1 Ohm? enota za upor.

Kirchhoffovi zakoni

a.) Za tokove:

$$I_1 + I_2 = I_3 + I_4$$

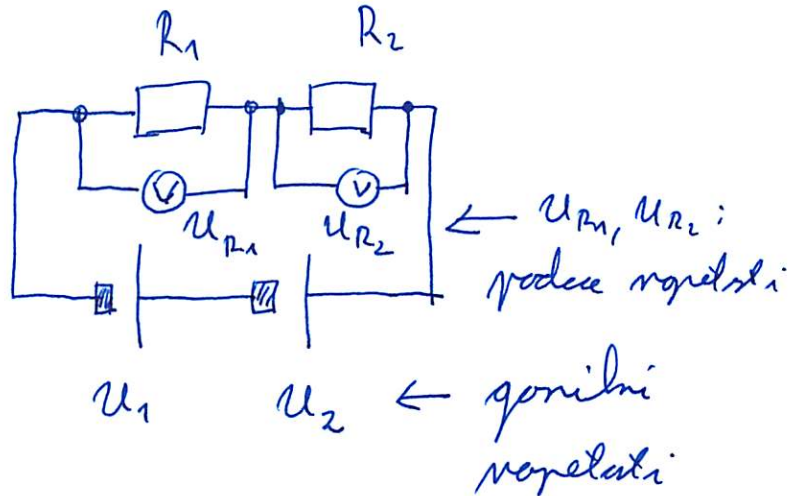


Uvoda tokov, ki tečejo v razvejitveni in evaha
 vstopi tokov, ki tečejo iz razvejitve!

b.) Za napetosti

$$\sum U_g = \sum U_p$$

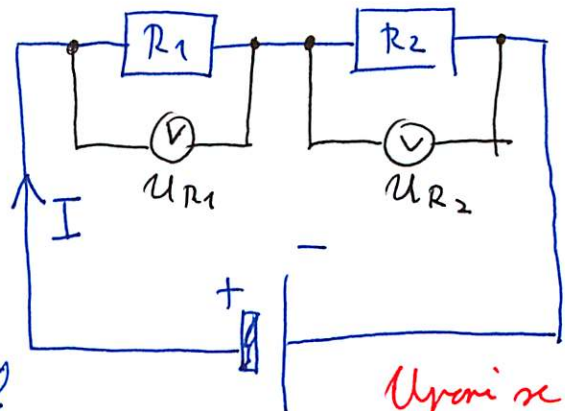
g: vsota gonilnih napetosti
 p: vsota padcev napetosti



Vežave uporov:

a.) Zaporedna vezava:

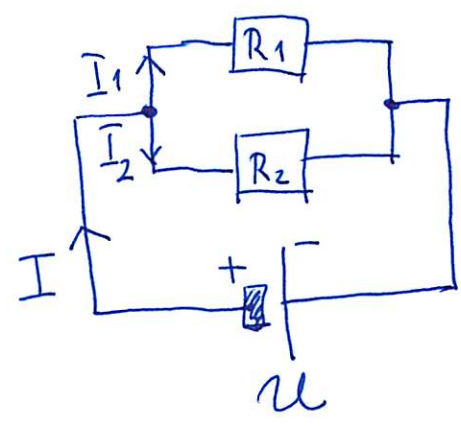
$U_{R1} + U_{R2} = U$ *indeksi!*
 Kirchhoffov zakon za napetosti
 $U_{R1} = I \cdot R_1; U_{R2} = I \cdot R_2$
 Ohmova zakon za posamezni upor?
 $I R_1 + I R_2 = U \Rightarrow (R_1 + R_2) I = U$



Upori se sestavljajo

$$R_m = R_1 + R_2$$

b.) Uzporedna vezava:



$I = I_1 + I_2$ Kirch. zah. za tokove

$U_{R1} = U$ in $U_{R2} = U$

Na vsakem uporu si vedec napetosti enak napetosti na izvira (gonilni napetosti)

$(I = I_1 + I_2)$

$U = U_{R1} = I_1 \cdot R_1$
 $U = U_{R2} = I_2 \cdot R_2$

$I = \frac{U}{R_1} + \frac{U}{R_2} = \left(\frac{1}{R_1} + \frac{1}{R_2} \right) U$
 $\frac{1}{R_m}$

$\frac{1}{R_m} = \frac{1}{R_1} + \frac{1}{R_2}$ Seštevanje se recipročne

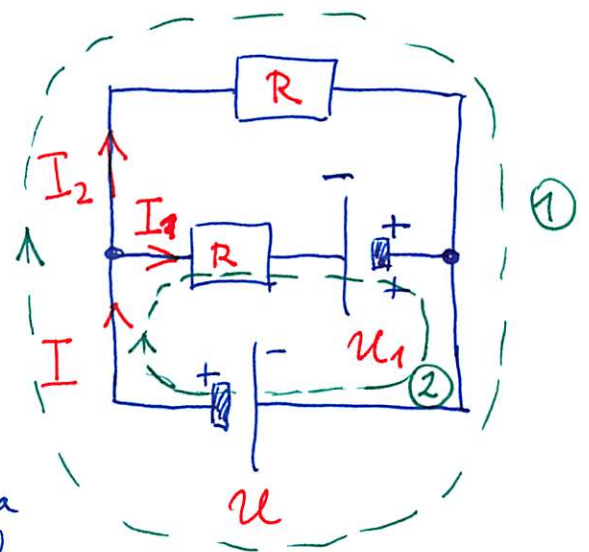
vnednosti uporov!

c.) Zled - primer vec gonilnih napetosti

$I = I_1 + I_2$

① $U = I_2 \cdot R$ Kirch. zahero za deso tokokrogca!

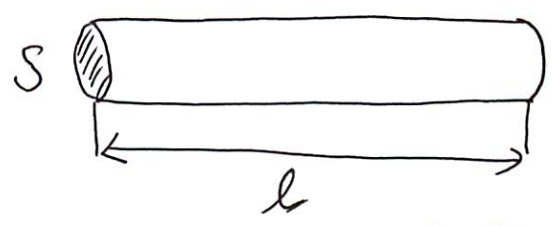
② $U + U_1 = I_1 R - I_2 R$



Podatki: U, U1 in R

$I = \frac{U + U_1}{R} + \frac{U}{R} = \frac{2U + U_1}{R}$

Specifični upor:



$$R = \xi \frac{l}{S}$$

	ξ	$[\Omega \text{mm}^2/\text{m}]$
Ag	0,015	
Cu	0,017	
Al	0,026	
Fe	0,1	

Električna moč:

$A = U \cdot e$ Delo elektrinske sile

$$dA = U de \quad | : dt$$

$$P = U I$$

$$\frac{dA}{dt} = P = U \frac{de}{dt} = U \cdot I \quad \uparrow \text{električna moč}$$

Pomembna: $P = U \cdot I$; U : podle napetosti na prevodniku ter I : el. tok skozi prevodnik!

Če upoštevamo $U = IR$, velja tudi:

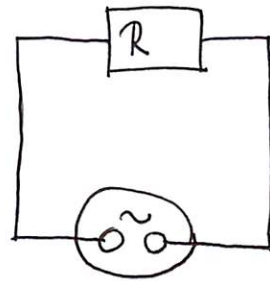
$$P = U I = \frac{U^2}{R} = I^2 \cdot R$$

Yzmerjena napetost in tok - Ohmski upor

$$U = I \cdot R$$

$$U_0 \sin \omega t = I \cdot R$$

$$I = \frac{U_0}{R} \sin \omega t = I_0 \sin \omega t$$



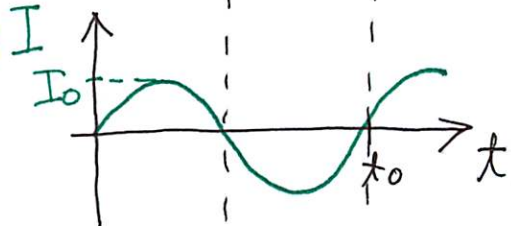
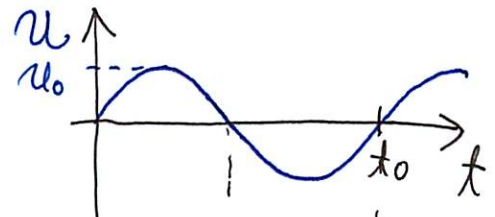
$$U = U_0 \sin \omega t$$

U in I nihata v fazi \rightarrow vzajna

Moč: $P = U \cdot I$

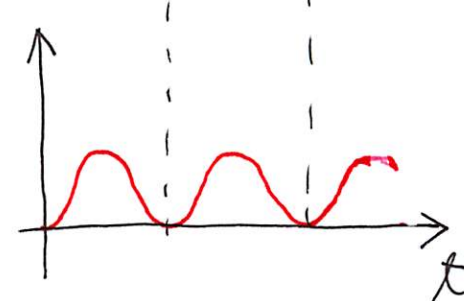
$$P(t) = U_0 I_0 \sin^2 \omega t$$

Moč nika v zero, vendar $P(t) > 0$



Sprosečna moč:

P



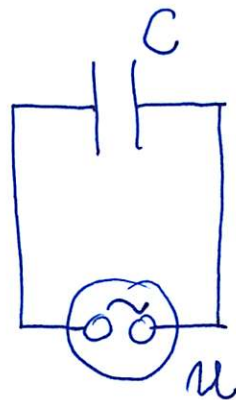
$$\bar{P} = \frac{1}{t_0} \int_0^{t_0} P(t) dt = \frac{U_0 I_0}{t_0} \int_0^{t_0} \sin^2(\omega t) dt =$$

$$= \frac{U_0 I_0}{t_0} \int_0^{t_0} \frac{1}{2} (1 - \cos 2\omega t) dt = \frac{U_0 I_0}{t_0} \left[\frac{t}{2} - \frac{\sin 2\omega t}{2 \cdot 2\omega} \right]_0^{t_0} =$$

$$\omega = \frac{2\pi}{t_0}$$

$$= \frac{U_0 I_0}{2}; \quad \bar{P} = \frac{U_0 I_0}{2} = \frac{U_0}{\sqrt{2}} \frac{I_0}{\sqrt{2}} = U_{ef} I_{ef}$$

Yzmenični tok skozi kondenzator



$$U = U_0 \sin \omega t$$

$$U = \frac{q}{C} \quad \text{oziroma} \quad q = C \cdot U$$

↑ napetost na izvira ↑ polje napetosti na kondenzatorju

$$q = C \cdot U \left| \frac{d}{dt} \right. ; \quad I = \frac{dq}{dt} !$$

$$I = C \frac{dU}{dt} = C \omega U_0 \cos \omega t = I_0 \cos \omega t$$

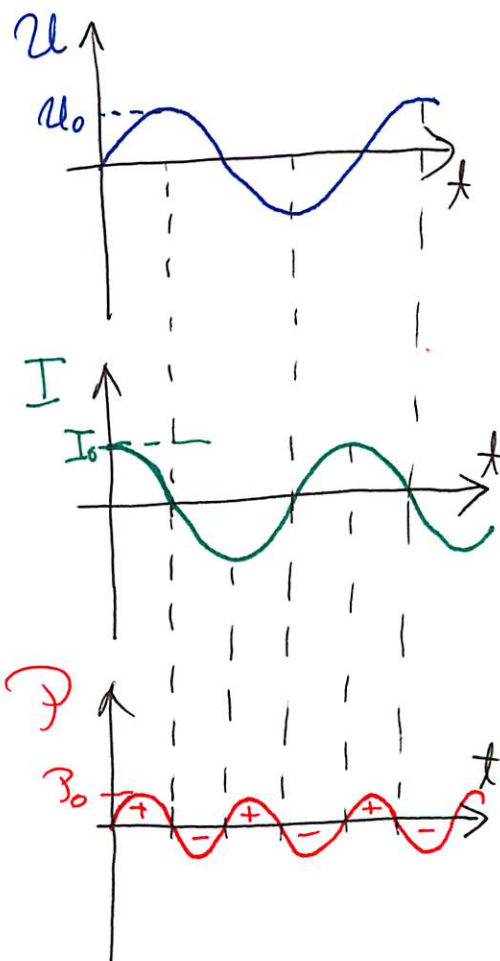
$$\Rightarrow U_0 = \frac{I_0}{C \omega} = R_C I_0 ; \quad \boxed{R_C = \frac{1}{C \omega}} : \text{upor } \text{kondenzatorja}$$

Moč, ki se troši na kondenzatorju:

$$P = U \cdot I = U_0 I_0 \sin \omega t \cos \omega t$$

$$P = \frac{U_0 I_0}{2} \sin 2\omega t$$

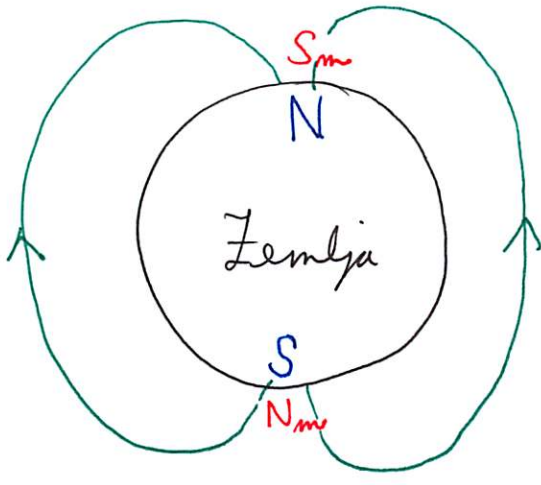
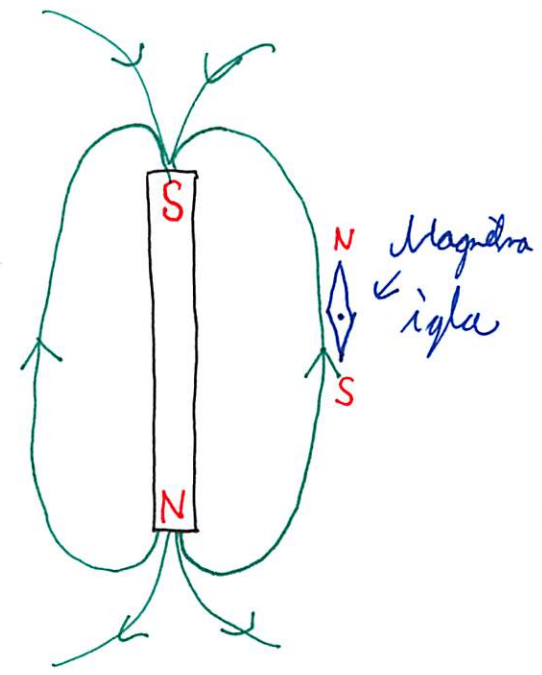
$$\overline{P} = \frac{1}{T_0} \int_0^{T_0} \frac{U_0 I_0}{2} \sin 2\omega t = 0 !$$



- železni gvilki

Magnetna polja

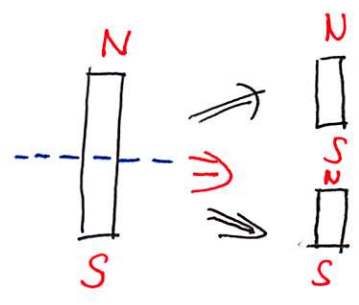
S južni magnetni pol
 N severni magnetni pol



Silnice magnetnega polja lahko "otipamo" s pomočjo magnetne igle

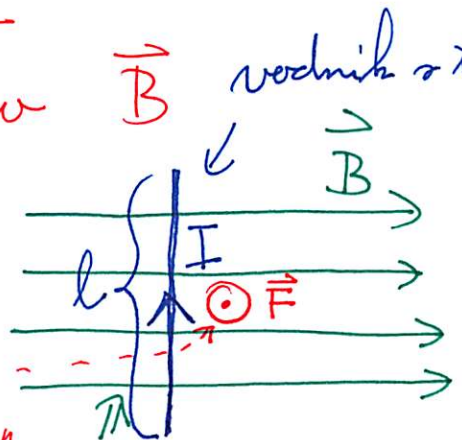
Magnetni magnoli ne obstajajo:

Če presečemo magnet na dva dela, dobimo dva magneta



Sila na vodnik v \vec{B} vodnik s tokom

$$\vec{F} = I \vec{l} \times \vec{B}$$

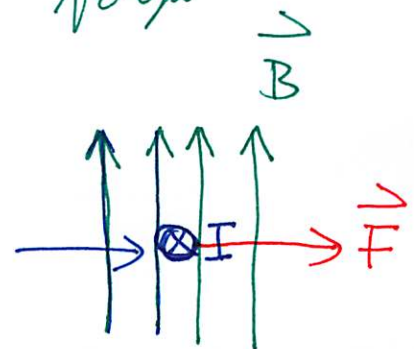


Sila krože "gor"! Silnice magnetnega polja

Enota?

$$B = \frac{F}{I l} \left[\frac{N \cdot m}{A \cdot m} = \frac{A \cdot V \cdot s}{A \cdot m^2} = \frac{V \cdot s}{m^2} = 1T \right]$$

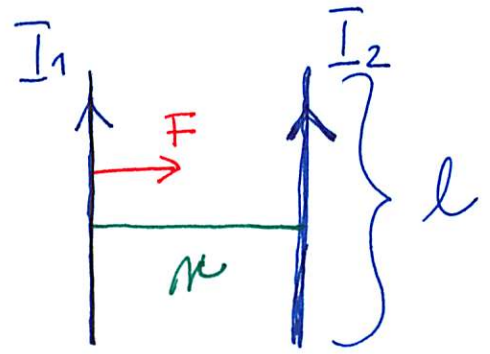
↑ gostota magnetnega polja se meri v Teslah!
 vodnik s tokom!



Sila med vodnikoma s tokom

tokom

$$F = \frac{\mu_0 l I_1 I_2}{2\pi r}$$

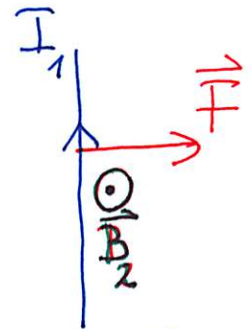


μ_0 : magnetna permeabilnost
vakuuma $\mu_0 = 4\pi \cdot 10^{-7} \frac{Vs}{Am}$

r : razdalja med paralelnima vodnikoma

I_1, I_2 : tokova, ki tečeta po vodnikih

l : dolžina



Povezava z enočlo: $F = I l B$

Vzemimo, da drugi vodnik povzroči magnetno polje B_2 na mestu prvega vodnika $\Rightarrow F = I_1 l B_2$

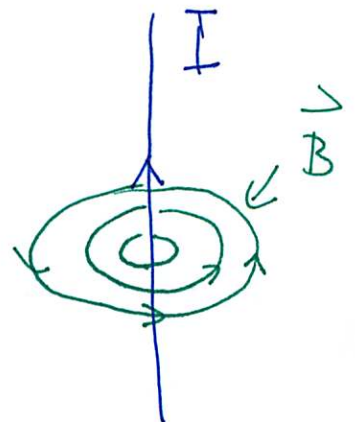
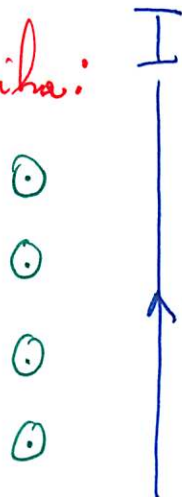
primerjajmo enočli: $F = \mu_0 \frac{l I_1 I_2}{2\pi r}$

ker $F = I_1 l B_2 \Rightarrow B_2 = \frac{\mu_0 I_2}{2\pi r}$

Ugotovitev posledica:

B v okolici neskončnega ravnega vodnika:

$$B = \frac{\mu_0 I}{2\pi r}$$

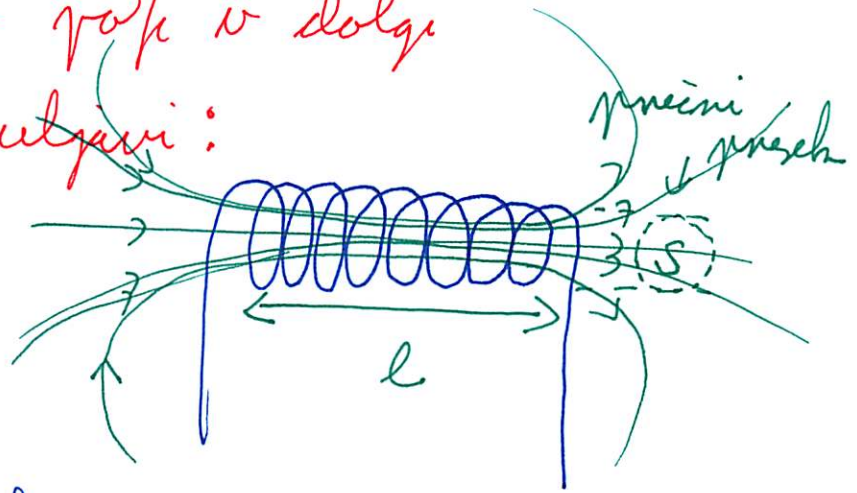


Risava se
projiciraji

Magnetno polje v dolgi

teljavi:

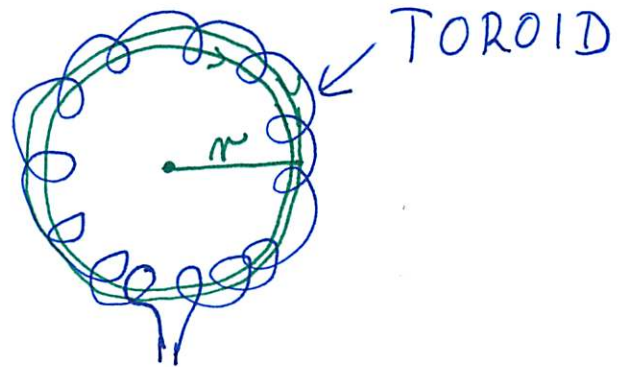
$$B = \frac{\mu_0 N I}{l}$$



Znotraj teljave je \vec{B} homogena in usmerjena vzdolž teljave.

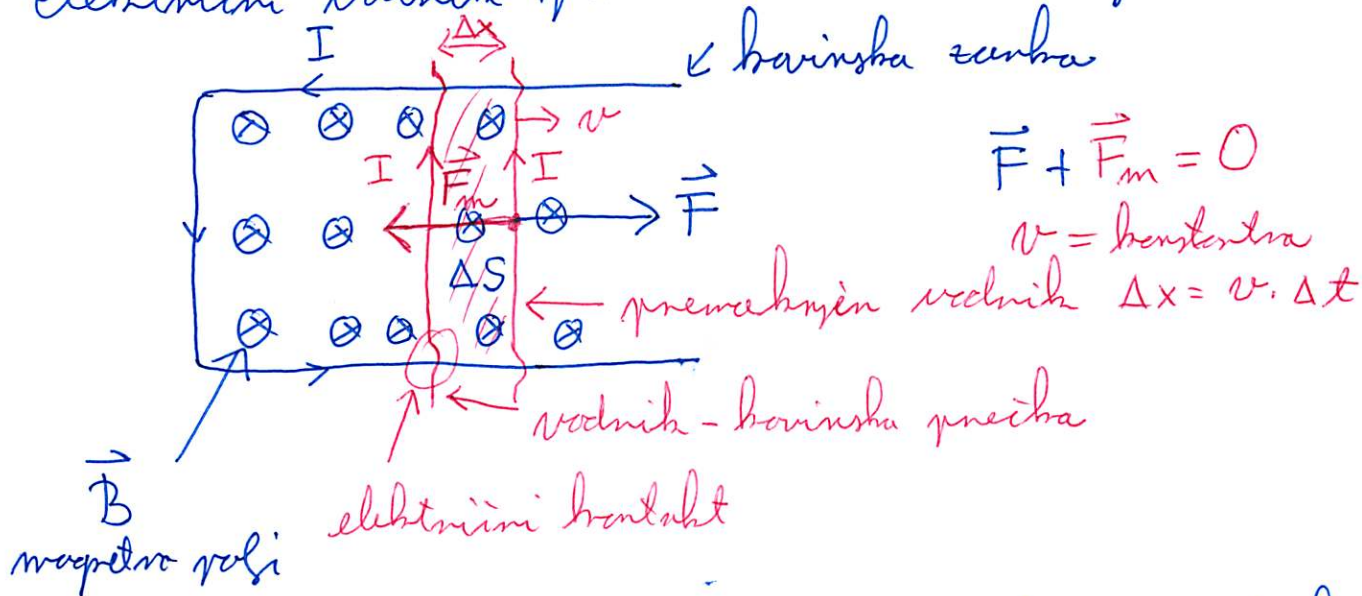
V toroidu:

$$B = \frac{\mu_0 N I}{2\pi r}$$



Indukcija

Električni vodnik premikamo skozi magnetno polje
 ← barinska zanka



Preko zanke ter prečke zavine teži električni tok

$$\vec{F}_m = I \vec{l} \times \vec{B} \quad F_m = F = I l B$$

Delo sile F v času Δt : $\Delta A = F \Delta x = F v \Delta t$

Električno delo: $\Delta A = \mathcal{P} \cdot \Delta t = \mathcal{U}_i I \Delta t$

\mathcal{U}_i : Inducirana napetost (napetost, ki vzganja el. tok na obkroju)

Uzročimo električno delo \mathcal{E} mehanskim delom sile F :

$$\mathcal{U}_i I \Delta t = F v \Delta t = \frac{I l B v \Delta t}{1} \quad \Delta \Phi_m$$

$$\mathcal{U}_i = l v B = \frac{l \Delta x}{\Delta t} B = \frac{\Delta S B}{\Delta t}$$

$\mathcal{U}_i = l v B$ ali $\mathcal{U}_i = \frac{\Delta \Phi_m}{\Delta t}$ } sprememba magnetnega pretoka po času.