

Heisenbergova interakcija za $S = 3/2$

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Naloga

Poišči lastne energije za operator Heisenbergove interakcije $H = -J\mathbf{S}_1 \cdot \mathbf{S}_2$, kjer ima vsak od spinov kvantno število velikosti vrtilne količine $S = 3/2$. Za (nekatera) lastna stanja sistema izračunaj pričakovane vrednosti in nedoločenosti komponent vrtilne količine v z smeri za posamezna spina.

Rešitev

V operatorju $H = -J\mathbf{S}_1 \cdot \mathbf{S}_2$ nastopa skalarni produkt obeh spinov. Skupni spin je $\mathbf{S} = \mathbf{S}_1 + \mathbf{S}_2$, njegov kvadrat pa se glasi $\mathbf{S}^2 = \mathbf{S}_1^2 + \mathbf{S}_2^2 + 2\mathbf{S}_1 \cdot \mathbf{S}_2$. Od tod lahko izrazimo $\mathbf{S}_1 \cdot \mathbf{S}_2 = \frac{1}{2}(\mathbf{S}^2 - \mathbf{S}_1^2 - \mathbf{S}_2^2)$ iz česar dobimo operator energije v obliki

$$H = -\frac{J}{2}(\mathbf{S}^2 - \mathbf{S}_1^2 - \mathbf{S}_2^2). \quad (1)$$

Lastna stanja za ta operator opišemo z velikostjo skupne vrtilne količine S in njeno projekcijo na os z , to je m . Vsako tako stanje lahko zapišemo kot vsoto različnih stanj obeh delcev, za katera mora veljati $|S_1 - S_2| \leq S \leq S_1 + S_2$ in $m = m_1 + m_2$, kjer sta S_1 ter S_2 velikosti obeh spinov in m_1 ter m_2 velikosti njunih projekcij na z os. Utežne koeficiente za vsako tako linearno kombinacijo stanj dobimo v Clebsch-Gordanovi tabeli koeficientov. S pomočjo razdelka $3/2 \times 3/2$ zapišemo vse možne lastne funkcije:

1. $|S = 3, m = 3\rangle = |m_1 = \frac{3}{2}, m_2 = \frac{3}{2}\rangle$
2. $|S = 3, m = 2\rangle = \sqrt{\frac{1}{2}}|m_1 = \frac{3}{2}, m_2 = \frac{1}{2}\rangle + \sqrt{\frac{1}{2}}|m_1 = \frac{1}{2}, m_2 = \frac{3}{2}\rangle$
3. $|S = 2, m = 2\rangle = \sqrt{\frac{1}{2}}|m_1 = \frac{3}{2}, m_2 = \frac{1}{2}\rangle - \sqrt{\frac{1}{2}}|m_1 = \frac{1}{2}, m_2 = \frac{3}{2}\rangle$
4. $|S = 3, m = 1\rangle = \sqrt{\frac{1}{5}}|m_1 = \frac{3}{2}, m_2 = -\frac{1}{2}\rangle + \sqrt{\frac{3}{5}}|m_1 = \frac{1}{2}, m_2 = \frac{1}{2}\rangle + \sqrt{\frac{1}{5}}|m_1 = -\frac{1}{2}, m_2 = \frac{3}{2}\rangle$
5. $|S = 2, m = 1\rangle = \sqrt{\frac{1}{2}}|m_1 = \frac{3}{2}, m_2 = -\frac{1}{2}\rangle - \sqrt{\frac{1}{5}}|m_1 = -\frac{1}{2}, m_2 = \frac{3}{2}\rangle$

6. $|S = 1, m = 1\rangle = \sqrt{\frac{3}{10}}|m_1 = \frac{3}{2}, m_2 = -\frac{1}{2}\rangle - \sqrt{\frac{2}{5}}|m_1 = \frac{1}{2}, m_2 = \frac{1}{2}\rangle + \sqrt{\frac{3}{10}}|m_1 = -\frac{1}{2}, m_2 = \frac{3}{2}\rangle$
7. $|S = 3, m = 0\rangle = \sqrt{\frac{1}{20}}|m_1 = \frac{3}{2}, m_2 = -\frac{3}{2}\rangle + \sqrt{\frac{9}{20}}|m_1 = \frac{1}{2}, m_2 = -\frac{1}{2}\rangle + \sqrt{\frac{9}{20}}|m_1 = -\frac{1}{2}, m_2 = \frac{1}{2}\rangle + \sqrt{\frac{1}{20}}|m_1 = -\frac{3}{2}, m_2 = \frac{3}{2}\rangle$
8. $|S = 2, m = 0\rangle = \sqrt{\frac{1}{4}}|m_1 = \frac{3}{2}, m_2 = -\frac{3}{2}\rangle + \sqrt{\frac{1}{4}}|m_1 = \frac{1}{2}, m_2 = -\frac{1}{2}\rangle - \sqrt{\frac{1}{4}}|m_1 = -\frac{1}{2}, m_2 = \frac{1}{2}\rangle - \sqrt{\frac{1}{4}}|m_1 = -\frac{3}{2}, m_2 = \frac{3}{2}\rangle$
9. $|S = 1, m = 0\rangle = \sqrt{\frac{9}{20}}|m_1 = \frac{3}{2}, m_2 = -\frac{3}{2}\rangle - \sqrt{\frac{1}{20}}|m_1 = \frac{1}{2}, m_2 = -\frac{1}{2}\rangle - \sqrt{\frac{1}{20}}|m_1 = -\frac{1}{2}, m_2 = \frac{1}{2}\rangle + \sqrt{\frac{9}{20}}|m_1 = -\frac{3}{2}, m_2 = \frac{3}{2}\rangle$
10. $|S = 0, m = 0\rangle = \sqrt{\frac{1}{4}}|m_1 = \frac{3}{2}, m_2 = -\frac{3}{2}\rangle - \sqrt{\frac{1}{4}}|m_1 = \frac{1}{2}, m_2 = -\frac{1}{2}\rangle + \sqrt{\frac{1}{4}}|m_1 = -\frac{1}{2}, m_2 = \frac{1}{2}\rangle - \sqrt{\frac{1}{4}}|m_1 = -\frac{3}{2}, m_2 = \frac{3}{2}\rangle$
11. $|S = 3, m = -1\rangle = \sqrt{\frac{1}{5}}|m_1 = \frac{1}{2}, m_2 = -\frac{3}{2}\rangle + \sqrt{\frac{3}{5}}|m_1 = -\frac{1}{2}, m_2 = -\frac{1}{2}\rangle + \sqrt{\frac{1}{5}}|m_1 = -\frac{3}{2}, m_2 = \frac{1}{2}\rangle$
12. $|S = 2, m = -1\rangle = \sqrt{\frac{1}{2}}|m_1 = \frac{1}{2}, m_2 = -\frac{3}{2}\rangle - \sqrt{\frac{1}{2}}|m_1 = -\frac{3}{2}, m_2 = \frac{1}{2}\rangle$
13. $|S = 1, m = -1\rangle = \sqrt{\frac{3}{10}}|m_1 = \frac{1}{2}, m_2 = -\frac{3}{2}\rangle - \sqrt{\frac{2}{5}}|m_1 = -\frac{1}{2}, m_2 = -\frac{1}{2}\rangle + \sqrt{\frac{3}{10}}|m_1 = -\frac{3}{2}, m_2 = \frac{1}{2}\rangle$
14. $|S = 3, m = -2\rangle = \sqrt{\frac{1}{2}}|m_1 = -\frac{1}{2}, m_2 = -\frac{3}{2}\rangle + \sqrt{\frac{1}{2}}|m_1 = -\frac{3}{2}, m_2 = -\frac{1}{2}\rangle$
15. $|S = 2, m = -2\rangle = \sqrt{\frac{1}{2}}|m_1 = -\frac{1}{2}, m_2 = -\frac{3}{2}\rangle - \sqrt{\frac{1}{2}}|m_1 = -\frac{3}{2}, m_2 = -\frac{1}{2}\rangle$
16. $|S = 3, m = -3\rangle = |m_1 = -\frac{3}{2}, m_2 = -\frac{3}{2}\rangle$

Iz enačbe (1) dobimo, če upoštevamo $S_1 = S_2 = \frac{3}{2}$, za lastne energije sledeč izraz:

$$H|S, m\rangle = -\frac{J\hbar^2}{2} \left(S(S+1) - \frac{3}{2} \left(\frac{3}{2} + 1 \right) - \frac{3}{2} \left(\frac{3}{2} + 1 \right) \right). \quad (2)$$

Iz enačbe je razvidno, da je energija posameznega lastnega stanja odvisna le od velikosti skupnega spina S . Za zapisana stanja dobimo lastne energije:

Energija	Lastna stanja (zaporedne številke)
$\frac{15}{2} \frac{J\hbar^2}{2}$	10.
$\frac{11}{2} \frac{J\hbar^2}{2}$	6., 9., 13.
$\frac{3}{2} \frac{J\hbar^2}{2}$	3., 5., 8., 12., 15.
$-\frac{9}{2} \frac{J\hbar^2}{2}$	1., 2., 4., 7., 11., 14., 16.

Na koncu izračunamo še nekaj pričakovanih vrednosti in nedoločenosti komponent vrtilne količine v z smeri za posamezna spina. Pričakovane vrednosti so

$$\langle \text{stanje} | S_{iz} | \text{stanje} \rangle, \quad \langle \text{stanje} | S_{iz}^2 | \text{stanje} \rangle, \quad i = 1, 2,$$

nedoločenosti pa

$$\delta S_{iz}^2 = \langle \text{stanje} | S_{iz}^2 | \text{stanje} \rangle - \langle \text{stanje} | S_{iz} | \text{stanje} \rangle^2.$$

Za stanje $S = 3, m = 3$ za prvi spin dobimo:

$$\langle S = 3, m = 3 | S_{1z} | S = 3, m = 3 \rangle = \left\langle m_1 = \frac{3}{2}, m_2 = \frac{3}{2} \left| S_{1z} \right| m_1 = \frac{3}{2}, m_2 = \frac{3}{2} \right\rangle = \frac{3\hbar}{2},$$

$$\langle S = 3, m = 3 | S_{1z}^2 | S = 3, m = 3 \rangle = \left\langle m_1 = \frac{3}{2}, m_2 = \frac{3}{2} \left| S_{1z}^2 \right| m_1 = \frac{3}{2}, m_2 = \frac{3}{2} \right\rangle = \frac{9\hbar^2}{4},$$

$$\delta S_{1z}^2 = \frac{9\hbar^2}{4} - \frac{9\hbar^2}{4} = 0.$$

Nedoločenost tega stanja je 0. Stanje $|m_1 = 3/2, m_2 = 3/2\rangle$ je namreč lastno stanje operatorja H .

Za stanje $S = 3, m = 2$ za drugi spin dobimo:

$$\begin{aligned} \langle S = 3, m = 2 | S_{2z} | S = 3, m = 2 \rangle &= \frac{1}{2} \left\langle m_1 = \frac{3}{2}, m_2 = \frac{1}{2} \left| S_{2z} \right| m_1 = \frac{3}{2}, m_2 = \frac{1}{2} \right\rangle + \\ &\frac{1}{2} \left\langle m_1 = \frac{3}{2}, m_2 = \frac{1}{2} \left| S_{2z} \right| m_1 = \frac{1}{2}, m_2 = \frac{3}{2} \right\rangle + \\ &\frac{1}{2} \left\langle m_1 = \frac{1}{2}, m_2 = \frac{3}{2} \left| S_{2z} \right| m_1 = \frac{3}{2}, m_2 = \frac{1}{2} \right\rangle + \\ &\frac{1}{2} \left\langle m_1 = \frac{1}{2}, m_2 = \frac{3}{2} \left| S_{2z} \right| m_1 = \frac{1}{2}, m_2 = \frac{3}{2} \right\rangle = \end{aligned}$$

$$= \frac{1}{2} \frac{\hbar}{2} + 0 + 0 + \frac{1}{2} \frac{3\hbar}{2} = \hbar \quad \text{in na enak način}$$

$$\langle S = 3, m = 2 | S_{2z}^2 | S = 3, m = 2 \rangle = \frac{1}{2} \frac{\hbar^2}{4} + 0 + 0 + \frac{1}{2} \frac{9\hbar^2}{4} = \frac{5\hbar^2}{4}, \quad \text{iz česar sledi}$$

$$\delta S_{2z}^2 = \frac{5\hbar^2}{4} - \hbar^2 = \frac{\hbar^2}{4}.$$