

Harmonski oscilator

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5. april 2004

Za harmonski oscilator velja:

$$\begin{aligned} H &= \frac{p^2}{2m} + \frac{1}{2}kx^2 \\ H &= \hbar\omega \left(a^\dagger a + \frac{1}{2} \right) \\ \omega &= \sqrt{\frac{k}{m}} \end{aligned} \tag{1}$$

Zapišimo še anihilacijski operator a in kreacijski operator a^\dagger :

$$a = \frac{1}{\sqrt{2}} \left(\frac{x}{x_0} + i \frac{p}{p_0} \right) \tag{2}$$

$$a^\dagger = \frac{1}{\sqrt{2}} \left(\frac{x}{x_0} - i \frac{p}{p_0} \right) \tag{3}$$

$$[a, a^\dagger] = 1 \tag{4}$$

Kjer sem označil:

$$x_0 = \sqrt{\frac{\hbar}{m\omega}} \quad \text{in} \quad p_0 = \frac{\hbar}{x_0}$$

$$E_n = \hbar\omega \left(n + \frac{1}{2} \right) \tag{5}$$

1 Zapis lastnih stanj

$$a^\dagger |n\rangle = \sqrt{n+1} |n+1\rangle \tag{6}$$

$$a |n\rangle = \sqrt{n} |n-1\rangle \tag{7}$$

$$|n\rangle = \frac{(a^\dagger)^n}{\sqrt{n!}} |0\rangle \tag{8}$$

2 Časovni razvoj

Na začetku imamo funkcijo:

$$|\Psi, 0\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \quad (9)$$

Hitro dobimo časovni razvoj te funkcije, saj imamo opravka z lastnimi stanji:

$$\begin{aligned} |\Psi, t\rangle &= \frac{1}{\sqrt{2}} \left(|0\rangle e^{-i\frac{E_0}{\hbar}t} + |1\rangle e^{-i\frac{E_1}{\hbar}t} \right) \\ &= \frac{1}{\sqrt{2}} \left(|0\rangle e^{-i\frac{\omega}{2}t} + |1\rangle e^{-i\frac{3\omega}{2}t} \right) \end{aligned} \quad (10)$$

Tu upoštevamo enačbo (5).

3 Pričakovane vrednosti

Poglejmo si, kako je s pričakovanimi vrednostmi za splošen $\langle A \rangle$:

$$\begin{aligned} \langle A \rangle &= \langle \Psi | A | \Psi \rangle = \left(\frac{1}{\sqrt{2}} \right)^2 \langle 0 | A | 0 \rangle + \left(\frac{1}{\sqrt{2}} \right)^2 \langle 1 | A | 1 \rangle + \\ &\quad + \frac{1}{2} \langle 0 | A | 1 \rangle e^{i\frac{\omega}{2}t} e^{-i\frac{3\omega}{2}t} + \frac{1}{2} \langle 1 | A | 0 \rangle e^{i\frac{3\omega}{2}t} e^{-i\frac{\omega}{2}t} \\ &= \frac{1}{2} \langle 0 | A | 0 \rangle + \frac{1}{2} \langle 1 | A | 1 \rangle + \text{Re} \left(\langle 1 | A | 0 \rangle e^{i\omega t} \right) \end{aligned} \quad (11)$$

3.1 Pričakovana vrednost koordinate

$$A = x$$

Hitro ugotovimo, da velja $\langle 0 | x | 0 \rangle = \langle 1 | x | 1 \rangle = 0$. Tako nam od enačbe (11) ostane le še:

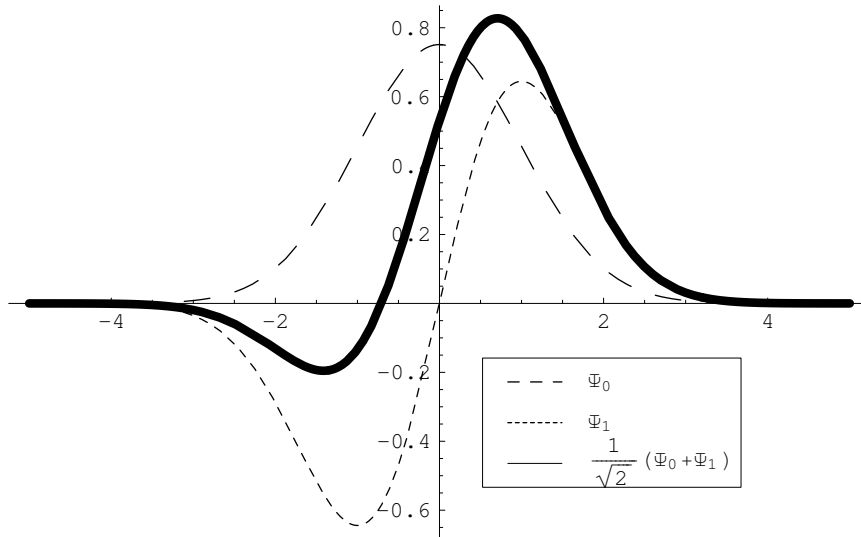
$$\langle x \rangle = \text{Re} \left(\langle 1 | x | 0 \rangle e^{i\omega t} \right) \quad (12)$$

Za lažje računanje x izrazimo z a in a^\dagger . Enačbi (2) in (3) seštejem in dobim:

$$\begin{aligned} a + a^\dagger &= \frac{2}{\sqrt{2}} x \\ x &= \frac{\sqrt{2}}{2} x_0 (a + a^\dagger) \end{aligned} \quad (13)$$

$$\begin{aligned}
\langle x \rangle &= \operatorname{Re} \left(\left\langle 1 \left| \frac{\sqrt{2}}{2} x_0 (a + a^\dagger) \right| 0 \right\rangle e^{i\omega t} \right) \\
&= \frac{\sqrt{2}x_0}{2} \operatorname{Re} \left[\left(\overbrace{\langle 1|a|0\rangle}^0 + \langle 1|a^\dagger|0\rangle \right) e^{i\omega t} \right] \\
&= \frac{\sqrt{2}x_0}{2} \operatorname{Re} \left(\langle 1|\sqrt{1}|1\rangle e^{i\omega t} \right) \\
\langle x \rangle &= \frac{\sqrt{2}}{2} x_0 \cos(\omega t)
\end{aligned} \tag{14}$$

Tu smo v drugi vrstici računa upoštevali ortonormiranost stanj: $\langle i|j\rangle = \delta_{i,j}$. Za boljšo predstavbo si pogledajmo, kako izgledata funkciji osnovnega $\Psi_0(x)$ in prvega vzbujenega $\Psi_1(x)$ stanja, ter kakšna je naša funkcija $\Psi(x) = \frac{1}{\sqrt{2}}(\Psi_0(x) + \Psi_1(x))$, ki je linearna kombinacija prvih dveh. Vse seveda za čas $t = 0$:



3.2 Pričakovana vrednost gibalne količine

$$A = p$$

Po enačbi (11) dobimo za $A = p$:

$$\begin{aligned}
\langle p \rangle &= \frac{1}{2} \left(\overbrace{\langle 0|p|0\rangle}^0 + \overbrace{\langle 1|p|1\rangle}^0 \right) + \operatorname{Re} \left(\langle 1|p|0\rangle e^{i\omega t} \right) \\
\langle p \rangle &= \operatorname{Re} \left(\langle 1|p|0\rangle e^{i\omega t} \right)
\end{aligned} \tag{15}$$

Podobno kot pri x , sedaj p izrazim z a in a^\dagger . Enačbi (2) in (3) odštejem in dobim:

$$\begin{aligned} a - a^\dagger &= \frac{2}{\sqrt{2}} i \frac{p}{p_0} \\ p &= i \frac{\sqrt{2}}{2} p_0 (a^\dagger - a) \end{aligned} \quad (16)$$

$$\begin{aligned} \langle p \rangle &= \text{Re} \left(\left\langle 1 \left| i \frac{\sqrt{2}}{2} p_0 (a^\dagger - a) \right| 0 \right\rangle e^{i\omega t} \right) \\ &= \frac{\sqrt{2}}{2} p_0 \text{Re} \left[\left\langle 1 \left| i \left(a^\dagger - \overbrace{a}^{\rightarrow 0} \right) \right| 0 \right\rangle + e^{i\omega t} \right] \\ &= \frac{\sqrt{2}}{2} p_0 \text{Re} \left(\langle 1 | \sqrt{1} | 1 \rangle i e^{i\omega t} \right) \\ \langle p \rangle &= -\frac{p_0}{\sqrt{2}} \sin(\omega t) \end{aligned} \quad (17)$$

3.3 Pričakovana vrednost kvadrata koordinate

$$A = x^2$$

Iz enačbe (13) dobim enačbo za x^2 :

$$x^2 = \frac{x_0^2}{2} (a^2 + a^{\dagger 2} + a a^\dagger + a^\dagger a) \quad (18)$$

$$\begin{aligned} \langle x^2 \rangle &= \frac{x_0^2}{2} \left(\frac{1}{2} \langle 0 | a a^\dagger | 0 \rangle + \frac{1}{2} \langle 1 | a a^\dagger + a^\dagger a | 1 \rangle + \text{Re}(0) \right) \\ &= \frac{x_0^2}{4} \left(\langle 0 | a | 1 \rangle + \sqrt{2} \langle 1 | a | 2 \rangle + \langle 1 | a^\dagger | 0 \rangle \right) \\ &= \frac{x_0^2}{4} (1 + \sqrt{2}\sqrt{2} + 1) \\ \langle x^2 \rangle &= x_0^2 \end{aligned} \quad (19)$$

4 Verjetnost

Sedaj bomo izračunali verjetnost $P_{x<0}$, da se delec nahaja na levi polovici potenciala.

$$P_{x<0} = \int_{-\infty}^0 |\Psi(x, t)|^2 dx \quad (20)$$

$$\Psi(x, t) = \frac{1}{\sqrt{2}} (\Psi_0(x, t) + \Psi_1(x, t)) \quad (21)$$

Zapišemo znani lastni fukciji harmonskega oscilatorja $\Psi_0(x, t)$ in $\Psi_1(x, t)$:

$$\Psi_0(x, t) = \frac{1}{\sqrt[4]{\pi x_0^2}} e^{-\frac{x^2}{2x_0^2}} e^{-\frac{i}{2}t\omega} \quad (22)$$

$$\Psi_1(x, t) = \frac{x\sqrt{2}}{x_0} \frac{1}{\sqrt[4]{\pi x_0^2}} e^{-\frac{x^2}{2x_0^2}} e^{-\frac{3i}{2}t\omega} \quad (23)$$

Iz enačbe (20) dobim:

$$P_{x<0} = \frac{1}{2} - \frac{\cos(\omega t)}{\sqrt{2}\pi} \quad (24)$$

Ta časovni potek verjetnosti, da se delec nahaja na levi strani potenciala, narišem na ustrezen graf, kjer pišem φ kot fazo $\varphi = t\frac{\omega}{2\pi}$:

